

Is our dynamical understanding of the circulation changes associated with the Antarctic ozone hole sensitive to the choice of reanalysis dataset?

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Abstract. This study quantifies differences among four widely used atmospheric reanalysis datasets (ERA5, JRA-55, MERRA-2, and CSFR) in their representation of the dynamical changes induced by springtime polar stratospheric ozone depletion in the Southern Hemisphere during 1980-2001, as part of S-RIP. The intercomparison is performed as part of the SPARC (Stratosphere–troposphere Processes and their Role in Climate) Reanalysis Intercomparison Project (S-RIP). The reanalyses are generally in good agreement in their representation of the expected strengthening of the lower stratospheric polar vortex during the austral spring-summer season, as well as the descent of anomalously strong winds to the surface during summer and the subsequent poleward displacement and intensification of the polar front jet. Differences in the trends in zonal wind are generally small compared to the mean trends. The exception is CSFR, which shows greater disagreement compared to the other three reanalysis datasets, with stronger westerly winds in the lower stratosphere in spring and a larger poleward displacement of the tropospheric westerly jet in summer. Our results suggest that that there is a high degree of consistency across the four reanalysis datasets in the representation of the dynamical changes associated with ozone depletion, which are examined by investigating the eddy heat and momentum fluxes and wave forcing. Nevertheless, there are larger differences in the wave forcing and eddy propagation changes compared to the similarities in the circulation trends. There is a large amount of disagreement in CFSR wave forcing / propagation trends compared to the other three reanalyses, while the best agreement is found between ERA5 and JRA-55. The underlying causes of these differences are consistent with the wind response being more constrained by the assimilation of observations compared to the wave forcing, which is more dependent on the model-based forecasts that can differ between reanalyses. Looking forward, these findings also give us confidence that reanalysis datasets can be used to assess changes associated with the ongoing recovery of stratospheric ozone.

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1 Introduction

Since the early 1980s, the polar stratosphere in the Southern Hemisphere (SH) has exhibited substantial cooling of up to 6-10 K during austral spring in response to the Antarctic ozone hole, driven by the reduction in radiative heating by stratospheric ozone (Randal and Wu, 1999; Thompson and Solomon, 2002). The cooling increases the meridional temperature gradient from the middle to high latitudes in the lower stratosphere, which is associated, via thermal wind balance, with a strengthening of the stratospheric polar vortex (Thompson and Solomon, 2002), and a subsequent delay in its springtime breakup (Keeble et al., 2014). In spring, the anomalously strong stratospheric winds propagate from around 10 hPa to the tropopause over the course of approximately a month, followed by a rapid descent through the troposphere in a few days (Thompson and Solomon, 2002).

The anomalous tropospheric circulation persists throughout austral summer and is characterized by a poleward shift of the extratropical jet stream, or polar front jet (Thompson and Solomon, 2002; Polvani et al., 2011). The tropospheric wind anomalies are fairly uniform throughout the troposphere (i.e., mostly barotropic), manifesting themselves at the surface as a shift in the midlatitude westerly winds, and are associated with a positive phase of the Southern Annual Mode (SAM) index (Thompson and Solomon, 2002; Marshall 2003; Arblaster and Meehl, 2006). The more positive SAM index has led to significant impacts on the regional climate of the extratropical SH (e.g., Gillett et al., 2006; Marshall et al., 2006, 2013; Orr et al., 2004, 2008; Van Lipzig et al., 2008; Thompson et al., 2011; Deb et al., 2018). Since the early 2000s, the stratospheric ozone has begun to show signs of recovery (Solomon et al., 2016), with the associated circulation trends slightly reversed or paused (Banerjee et al., 2020).

The dynamical basis of the polar front jet (or the SAM index) involves positive feedbacks between the anomalous westerlies and synoptic-scale eddy fluxes of momentum and heat. The stronger westerlies are accompanied by enhanced transient baroclinic eddy generation, which tend to propagate upward and equatorward from their latitudes of generation, resulting in a flux of momentum into the jet (convergence) that plays a major role in maintaining its persistence and mid-latitude variability (Robinson, 1996, 2000; Lorenz and Hartmann, 2001; Hartmann and Lo, 1998; Gerber and Vallis, 2007).

The stratospheric polar vortex is strongly influenced by planetary-scale waves propagating up from the troposphere (Christiansen, 1999; Plumb, 2010), which are associated with eddy heat fluxes and play an important role in transferring heat from low to high latitudes. Heating perturbations in the stratosphere, which alter the meridional temperature gradient (and via thermal wind balance the vertical shear of background winds), have been shown to modulate the upward propagation of planetary waves, with changes at the tropopause key to controlling the amount of wave activity transferred from the troposphere into the stratosphere (Matsuno, 1970; Chen and Robinson, 1992; Scott and Polvani, 2006; Martineau et al., 2018a). Furthermore, the attendant stratospheric circulation anomalies can propagate downwards to the tropopause, with the stratospheric forcing subsequently able to alter tropospheric annular modes by changing the synoptic-scale eddy feedbacks

that maintain variations in the polar front jet (Christiansen, 1999, 2001; Kodera and Kuroda, 2002; Limpasuvan et al., 2004; Song and Robinson, 2004; Smith and Scott, 2016). Consequently, a number of studies have suggested that the circulation changes induced by the ozone hole involve alterations to these dynamical processes, although the exact mechanisms remains
60 uncertain (e.g., Hartmann et al., 2000; Chen and Held, 2007; McLandress et al., 2010, 2011; Harnik et al., 2011, Shaw et al., 2011; Orr et al., 2012, 2013; Hu et al., 2015).

Orr et al. (2012) performed a model-based study focused on a momentum budget analysis within the Transformed Eulerian Mean (TEM) framework. It was used to test the hypothesis that the circulation changes associated with the ozone hole were caused by changes to wave forcing and are supported by strong dynamical feedbacks. They found that ozone depletion causes
65 an initial (radiative) strengthening of the lower-stratospheric winds, which results in a reduction of upward-propagating planetary waves from the troposphere into the stratosphere. This reduction of wave activity in the lower-stratosphere causes a decrease in the wave-driven deceleration of the polar vortex, resulting in its acceleration. This initiates a positive feedback process in which fewer planetary waves propagate up from the troposphere, further drawing the decrease in wave-driven deceleration and associated strengthened winds downwards to the tropopause. In addition, a confinement of planetary-waves
70 in the high-latitude troposphere is necessary for the coupling of the stratospheric and tropospheric changes. The positive feedback processes involving increases in low-level baroclinicity and the subsequent generation of baroclinic activity results in changes to the synoptic-scale wave fluxes of heat and momentum, which are important for the poleward displacement of the polar front jet in the troposphere. Finally, increased upward fluxes of planetary wave activity from the troposphere into the lower stratosphere at high latitudes occurs due to the delay in the breakup of the stratospheric vortex, resulting in stronger
75 wave-driven deceleration in summer.

Atmospheric reanalysis datasets combine observations and temporally unchanged weather forecast model information in an optimal way to construct a ‘best’ estimate of the state of the atmosphere. They have previously been used to investigate the circulation trends in the stratosphere and troposphere that occur in response to the ozone hole (e.g. Chen and Held, 2007; Harnik et al., 2011; Shaw et al., 2011; Orr et al., 2012, 2013; Banerjee et al., 2020). These studies tend to be based on a single
80 reanalysis dataset, despite many others being available (Fujiwara et al., 2017). Reanalysis systems assimilate both conventional data (e.g. radiosonde profiles, surface measurements, and aircraft measurements) and satellite data (e.g. infrared and microwave radiances). The availability of satellite data has increased substantially since the appearance of the ozone hole and contributed to major improvements in accuracy (Marshall, 2003; Sterl, 2004), as prior to the “satellite era” reanalyses are considered unreliable in the high latitudes of the SH due to the sparseness of conventional observations. As reanalysis datasets
85 largely use the same available input observations, differences in the technical details of the reanalysis systems means that they may give different results for the same diagnostics (Fujiwara et al., 2017). Key differences in current reanalysis systems include the data assimilation strategy, such as three- and four-dimensional variational (3D-VAR and 4D-VAR) approaches, as well as 3D-FGAT (First Guess at Appropriate Time). 4D-VAR makes better use of observations than either 3D-VAR or 3D-FGAT,

resulting in substantial improvements (Fujiwara et al., 2017). Differences in the forecast models used are also important, as they have their own biases throughout the atmosphere. Therefore, reanalysis datasets do not necessarily agree on how the SH circulation responds to the ozone hole, possibly making the results reanalysis dependent. This is perhaps especially an issue in the stratosphere, as compared to the troposphere this region is characterised by smaller volumes of observational data available for assimilation and larger biases in observational data (Fujiwara et al., 2017), implying a greater reliance on the performance of the forecast model and its representation of dynamical processes (e.g., Orr et al., 2010). The representation of the underlying dynamics in reanalyses is therefore an additional concern, which has not been examined for the SH despite showing nonnegligible differences for some diagnostics in the Northern Hemisphere (e.g., Bengtsson et al., 2006; Lu et al., 2015; Martineau et al., 2016, 2018b; Chemke and Polvani, 2020).

The primary aim of this study is to compare trends in the SH circulation over the 1980 to 2001 period associated with the ozone hole in four widely-used reanalyses, and to analyse their connection to changes in various dynamical quantities to establish whether they consistently support the proposed mechanisms associated with the ozone hole. The four reanalysis datasets examined are JRA-55 (Kobayashi et al. 2015), MERRA-2 (Gelaro et al., 2017), CFSR (Saha et al., 2010, 2014), and ERA5 (Hersbach et al., 2020).

2 Data and methods

Details of the four reanalysis systems examined are given in Table 1. See also Fujiwara et al. (2017) for a summary of each of the reanalysis systems. Key differences are that both ERA5 and JRA-55 employ a 4D-VAR (Kobayashi et al. 2015; Hersbach et al., 2020) scheme, while CSFR employs a 3D-VAR scheme (Saha et al., 2014). MERRA-2 employs a 3D-FGAT scheme (Lawless, 2010), which is regarded as an intermediate step between 3D-VAR and 4D-VAR. There is a considerable difference in the release date of the forecast model used by each system, which is relevant as the models will have improved over time. ERA5 and MERRA-2 use considerably more recent model versions (year 2016 and 2015, respectively) compared to JRA-55 and CSFR (year 2009 and 2007, respectively). The four systems have a broadly similar horizontal grid spacing of 0.5° or better.

The data used in this study are described by Martineau et al. (2018c) and were produced as part of the Stratosphere-troposphere Processes And their Role in Climate (SPARC) Reanalysis Intercomparison Project (S-RIP) to facilitate the comparison of reanalysis datasets. They include zonally averaged atmospheric diagnostics of basic dynamical variables and more advanced wave forcing quantities computed using the four reanalyses datasets examined. The variables are provided every six hours and prepared using a common $2.5^\circ \times 2.5^\circ$ grid and standard pressure levels. For this investigation a subset of the data was retrieved, comprising the period from 1980 to 2001 and 15 pressure levels (1000, 925, 850, 700, 600, 500, 400, 300, 250, 200, 150, 100, 70, 50, 30 hPa). Note that no data is provided for MERRA-2 in the range 1000-500 hPa since unlike the other reanalyses it

does not provide data extrapolated below the surface. Additionally, ERA5 exhibits a pronounced cold bias in lower
120 stratospheric temperature from 2000 to 2006 due to the use of inappropriate background error covariances (Simmons et al.,
2020). This issue was fixed in a new set of ERA5 reanalysis from 2000 to 2006, termed ERA5.1 (Simmons et al., 2020), which
we used instead of ERA5 for this period (hereinafter this combined dataset is referred to as ERA5 for simplicity).

The key variables examined in this study are the zonally averaged zonal wind \bar{u} , the eddy momentum flux $\overline{u'v'}$, and the eddy
heat flux $\overline{v'T'}$. Here T is the temperature, u the zonal wind, v the meridional wind, overbars denote zonal means, and primes
125 denote deviations from the zonal mean. Additionally, the quasi-geostrophic form of the TEM momentum equation is used to
diagnose the wave forcing on the zonally averaged zonal wind (Edmon et al., 1980). This is expressed as

$$\frac{\partial \bar{u}}{\partial t} = f\bar{v}^* + \frac{1}{a \cos \phi} \nabla \cdot F^{QG} + \bar{\epsilon} \quad (1)$$

where F^{QG} is the quasi-geostrophic Eliassen-Palm (EP) flux, which takes the form

$$F^{QG} = \{F_{\phi}^{QG}, F_p^{QG}\} = a \cos \phi \left\{ -\overline{u'v'}, \frac{\overline{v'\theta'}}{\partial \theta / \partial p} f \right\} \quad (2)$$

130 with the wave forcing represented by the EP flux divergence (EPFD) term, which is the second term on the right-hand-side of
Eq. (1). The first term on the right-hand-side of Eq. (1) is the Coriolis torque. Here f is the Coriolis frequency, \bar{v}^* is the residual
meridional circulation, a is the mean radius of the Earth, ϕ is the latitude, ϵ represents any residual tendencies (unresolved
waves, diffusion, ageostrophic effects), θ is the potential temperature, and p is pressure (Martineau et al., 2018c).

We examine the momentum flux and heat flux instead of the EP flux components $(F_{\phi}^{QG}, F_p^{QG})$ as the latter requires the vertical
135 derivative of temperature or static stability, resulting in noisy wave driving and EP fluxes (Lu et al., 2014). The eddy heat
fluxes play a key role in the vertical component of EP flux (F_p^{QG}) , which is a measure of the upward fluxes of Rossby wave
activity (Edmon et al., 1980). In the SH, positive (negative) anomalies of the eddy heat flux indicate reduced (enhanced)
poleward heat transfer, while positive (negative) anomalies of the eddy momentum flux indicate reduced (enhanced) poleward
momentum transfer.

140 The results compare linear trends over a 22-year period from 1980 to 2001 for the four reanalyses, focusing on austral spring
(September-October-November; SON) and summer (December-January-February; DJF). The 1980 to 2001 period was chosen
because the time-series of the ozone mass deficit measure of Huck et al. (2007) revealed that the ozone hole first emerged
around 1980, with its size steadily increasing until around 2000/2001. So selecting the 1980 to 2001 period maximises trends
in circulation and related dynamical quantities (Banerjee et al., 2020). It also provides a clean case study for reanalysis data

145 inter-comparison in terms of atmospheric trends and the associated dynamical connection between the troposphere and the stratosphere in the SH.

Results based on meridionally-averaged values of the zonal wind, eddy momentum flux $\overline{u'v'}$, eddy heat flux $\overline{v'T'}$, and EP flux divergence (hereinafter EPFD) are areal weighted using the cosine of latitude. To better compare differences between the reanalyses, ERA5 is chosen as a reference dataset and differences between it and MERRA-2, JRA-55 and CFSR are calculated.

150 The choice of ERA5 as a reference is somewhat arbitrary, in that we have no a priori expectation that it is closer to the truth. It is, however, the most recently developed reanalysis (Table 1). Furthermore, the vertically integrated momentum flux and heat flux are also computed (Held and Philipps, 1993) for all waves, as well as planetary (zonal wavenumbers 1-3) and synoptic (zonal wavenumbers 4 and higher) waves. This is to investigate differences in the propagation of synoptic waves in the troposphere and planetary waves in the stratosphere and their contributions to the total wave fluxes. Finally, the trend in

155 the final warming date of the Antarctic polar vortex was calculated using the method of Black and McDaniel (2007), which identified the final warming date as the final time that the zonally averaged wind at 60°S and 50 hPa falls below 10 m s⁻¹.

3 Results

Since SH winds have undergone such large changes in response to the ozone hole, we first compare the trends in the zonal wind among the reanalysis datasets. Figure 1 shows height (in pressure coordinates) versus latitude cross-sections of DJF

160 trends in zonally averaged zonal wind for the period 1980 to 2001 for the four reanalyses. All four reanalyses show the expected stronger westerly winds in both the stratosphere and troposphere. For ERA5, the strongest increases in both the troposphere (up to 2 m s⁻¹ dec⁻¹) and stratosphere (2.5 m s⁻¹ dec⁻¹) are confined to a relatively narrow latitudinal band of 55-65°S, although in the stratosphere the enhanced westerly winds expand equatorward to 30°S and poleward to 80°S, which is consistent with an overall strengthening of the climatological polar vortex. In the troposphere the strengthened westerly winds form part of a

165 dipole pattern, with easterly winds at around 40°S, which is consistent with a poleward shift of the polar front jet. The results for ERA5, JRA-55, and MERRA-2 are largely in good agreement. For CFSR, the peak wind increase in the stratosphere exceeds 3.0 m s⁻¹ dec⁻¹, and in both the stratosphere and troposphere the region of maximum increase in the westerlies is located in the range 60-70°S. This is further poleward and larger in magnitude in comparison with ERA5, resulting in the positive differences at 60-70°S and negative differences at 50-60°S when compared to ERA5. Note that CFSR also disagrees

170 with the other three reanalyses in terms of the corresponding (DJF) trends in temperature, evident by enhanced warming below 300 hPa (by ~0.4 K dec⁻¹) and cooling between 300 and 100 hPa (by ~-1 K dec⁻¹) relative to ERA5, resulting in comparative weakening of the stability near the tropopause (Figure A1).

Following Thompson and Solomon (2002), Figure 2 shows the corresponding time-height cross-sections of the trends in zonally averaged zonal wind (averaged over 50-70°S) from September to February. The expected strengthening of the winds

175 and their descent from the lower stratosphere into the troposphere is apparent in all four reanalyses. For ERA5, the trends in zonal wind can be separated into four stages: i) stronger westerly winds appearing in the lower stratosphere in September, ii) continued strengthening of the lower stratospheric winds from September to early December (peaking at $4 \text{ m s}^{-1} \text{ dec}^{-1}$) and descent to the tropopause, iii) weakening of the anomalously strong westerly winds in the lower stratosphere from December to January and descent of the winds to the surface, and iv) a continued weakening of the anomalously strong stratospheric winds from December to February, consistent with a delayed breakup of the vortex in summer. According to Orr et al. (2012), these four stages refer respectively to the ‘onset’, ‘growth’, ‘decline’, and ‘decay’ stages of the lifecycle of the zonal wind response to the ozone hole.

The results for ERA5, JRA-55, and MERRA-2 are again largely in good agreement (with differences not exceeding $\pm 0.6 \text{ m s}^{-1} \text{ dec}^{-1}$). The largest differences among the reanalyses are again associated with CFSR, which shows much stronger stratospheric winds than ERA5 between September and November (i.e., ‘onset’ and ‘growth’ stages), suggesting the initial strengthening of the winds occurs earlier in CFSR. Furthermore, the four reanalyses generally show a similar delay in the breakup of the polar vortex. The final warming date for all reanalyses occurs around 0.9 days later per year or around 19 days later over the period 1980 to 2001 (not shown).

190 In ERA5 the corresponding time-height cross-section of trends in zonally averaged temperature (Figure A1) demonstrates that the stratospheric cooling associated with the ozone hole lasts from October to January, with a peak of -4 K dec^{-1} in November. This agrees with radiosonde observations from Antarctica (Thomson and Solomon, 2002), and is also in agreement with MERRA-2 and JRA-55 results. However, CFSR again contrasts with the other three reanalyses in terms of the temperature trend, evident by both an earlier onset to the cooling (beginning from September) and enhanced cooling between 300 and 100 hPa (by $\sim -1 \text{ K dec}^{-1}$) throughout September to February.

195 To further investigate the response of the tropospheric polar front jet during DJF, Figure 3 shows the latitudinal profile of the trend and climatology of the 500 hPa zonally averaged zonal wind for the four reanalyses. The climatologies are nearly identical except poleward of $\sim 70^\circ\text{S}$ and show that the peak winds associated with the jet occur around 50°S . The lack of agreement poleward of $\sim 70^\circ\text{S}$ may be due to a lack of observations over the continent and/or the increase in uncertainty of zonal mean quantities near the pole, an effect of spherical geometry. Positive trends ($\sim 1.5 \text{ m s}^{-1} \text{ dec}^{-1}$) are found on the poleward flank of the jet while negative trends ($\sim -0.8 \text{ m s}^{-1} \text{ dec}^{-1}$) occur at $\sim 38^\circ\text{S}$, which is consistent with the results of Figure 1, i.e. a strengthened and poleward shift of the polar front jet in the troposphere. In comparison with other reanalyses, there is a clear poleward shift of $\sim 4^\circ$ for the CFSR trend, which is also consistent with the stronger poleward shift in the jet shown in Figure 1. The good agreement between the climatological results suggests that the differences in the trends are not due to biases/differences in the climatological strength or location of the tropospheric westerly jet.

To study the spread among the four reanalyses in terms of wave driving, Figure 4 (a,d,g,j) shows time-height cross-sections of the trend in EPFD (averaged over 40-80°S) from September to February. For ERA5, in the lower stratosphere the EPFD shows a positive trend during November (i.e., weaker wave drag, coinciding with the ‘growth’ stage and the peak increase in stratospheric winds), followed by a negative trend during DJF (i.e. stronger wave drag, coinciding with the ‘decay’ and ‘decline’ stages and a weakening of the strengthened vortex and its delayed breakup). This is in dynamical agreement with the temporal evolution of the zonal wind trends in Figure 2 but does not necessarily indicate causality. The total zonal wind acceleration (in the absence of e.g. unresolved small-scale forcing) is largely a balance between the Coriolis torque on the residual meridional circulation and the wave drag on these time scales (Eq. 1). For September and October, the trend in lower stratospheric EPFD is largely negligible, suggesting that the circulation response during this time is primarily radiatively controlled. Both positive and negative trends in EPFD descend from 30 hPa to 300 hPa, indicating a downward influence from the stratosphere. In the lower stratosphere the trend in EPFD shows little difference among the four reanalyses.

Orr et al. (2012) also describe a switch from weaker (in November) to stronger (in DJF) wave drag in response to the ozone hole. They emphasize two factors, (i) a positive feedback process whereby an initial strengthening of the polar vortex winds in response to radiative cooling (during the ‘onset’ phase) plays an important role in conditioning the polar vortex so that that fewer planetary-scale waves can propagate up from the troposphere, resulting in weaker wave drag (during the ‘growth’ phase): this is consistent with the conclusion of Chen and Robinson (1992) that increased vertical wind shear at the tropopause is key to reducing the propagation of planetary waves into the stratosphere. And (ii), a negative feedback process whereby the prolonged existence of the westerly winds due to the delay in the breakup of the stratospheric vortex permits increased upward wave propagation into the stratosphere, resulting in stronger wave drag (during the ‘decline’ and ‘decay’ stages): this is consistent with a larger “cavity” of positive refractive index in this region (wave activity tends to propagate towards more positive refractive index values).

In the troposphere, EPFD shows bands of negative (positive) trends in the upper (middle) troposphere for ERA5 from September through to February (cf. Figure 4a). The agreement among the four reanalyses is poor, with the discrepancies relative to ERA5 marked by alternating negative and positive horizontal stripes, which can be greater in amplitude than the mean trends, and are most prominent for CFSR (e.g., during October). However, the rather large spread in the tropospheric EPFD trends (Figure 4 (a,d,g,j)) are accompanied by relatively small differences in the tropospheric wind trends (Figure 2). There is also no evidence of vertically alternating differences in the wind trend.

These results suggest that in the troposphere the resolved EPFD trend is not directly linked to the trends in the zonal wind; the latter is more linked to direct observation, while the former is more forecast model dependent. In addition, the tropospheric circulation is relatively more constrained by observational input in comparison to the stratospheric circulation (Martineau et

al., 2016). Lu et al. (2014) found similar alternating stripes in the EPFD when they compared wave driving between ERA-Interim and ERA-40 reanalyses. They showed that one of the main contributors to the EPFD differences was the vertical derivative of the temperature. Note that interpolation from model levels to standard pressure surfaces could also play a role in discrepancies of the EPFD term, as derivatives are very sensitive to interpolation. Differences in trends in the upward component of the EP flux (Eq. 2), which also includes the vertical derivative of temperature, also characterized by alternating negative and positive horizontal stripes (not shown).

3.2 A dynamical analysis of trends: Eddy heat and momentum fluxes

Figure 4 (b,e,h,k) shows the time-height cross-sections of the trend in zonally averaged eddy heat flux $v'T'$ for the four reanalyses. The ERA5 results show a region of positive trend in the lower stratosphere in November indicating reduced poleward eddy heat flux / upward wave activity from the troposphere into the stratosphere, which corresponds to the positive trends in EPFD, i.e., reduced EP flux convergence. Comparison with Figure 2 shows that this period is also contemporaneous with the descent of the anomalously strong westerly winds / increased vertical wind shear to the tropopause. For DJF, the ERA5 results show a negative trend in the lower stratosphere signifying enhanced poleward eddy heat flux / upward propagating wave activity into the stratosphere, which corresponds to negative trends in EPFD, i.e., increased EP flux convergence. For September and October, the trend in lower stratospheric eddy heat flux is much smaller and noisier. This corresponds to the switch from weaker (in November, during the ‘growth’ stage) to stronger (in DJF, during the ‘decline’ and ‘decay’ stages) wave activity propagating into the lower stratosphere described by Orr et al. (2012). The other reanalyses exhibit minor differences compared to ERA5, except for CSFR, which exhibits a stronger negative trend of the eddy heat flux in DJF (and September and October) and a weaker positive trend in November. Additionally, in ERA5 the region of positive trend in heat flux in November appears to start from around the tropopause and extends upward quickly in time, while this effect is less apparent or more barotropic in the other three reanalyses. Negligible trends in the heat flux can be detected in the troposphere, confirming that changes in the upward propagating waves are confined in the stratosphere (Orr et al., 2012).

Figure 4 (c,f,i,l) shows the time-height section of the trend in zonally averaged eddy momentum fluxes $\overline{u'v'}$ for the four reanalyses. For ERA5, a negative trend is found to dominate the lower stratosphere from October to November, indicating enhanced poleward momentum transfer. Hartmann et al. (2000) argued that the enhanced vortex winds / vertical shear in the lower stratosphere associated with the ozone hole cause enhanced equatorward propagation of planetary waves, thus more negative $\overline{u'v'}$ in the SH. For the other three reanalyses, the negative stratospheric trend is stronger compared to ERA5, especially in CFSR (consistent with its stronger vortex winds from September to December (Figure 2), which favors increased equatorward wave propagation).

In the troposphere, in ERA5 the trend in eddy momentum flux is marked by persistent negative values from December to February, indicating enhanced poleward momentum transfer. This occurs at the same time as the poleward displacement of

the polar front jet and anomalously strong westerlies in the troposphere (Figures 1 and 2). This negative trend in eddy momentum flux in the troposphere is evident for all four reanalyses products, although JRA-55, MERRA-2, and CFSR have weaker trends than ERA5. Orr et al. (2012) similarly describe strengthened equatorward synoptic-scale wave propagation in the troposphere in response to the ozone hole during the ‘decline’ and ‘decay’ stages. They show that this coincides with enhanced baroclinity at the surface (i.e., an increase in upward synoptic-scale waves) at the same latitude as the strengthened polar front jet. This suggests that the circulation trends are the result of the interactions between the zonal-mean flow and the eddies, which maintain anomalies in the polar front jet / tropospheric annular mode. The fluxes of momentum into the jet (convergence) balances anomalous surface wind stress associated with the shift (see also Hartmann et al., 2000).

The analysis in the next two sub-sections further explores the differences in the trends in eddy heat and momentum fluxes for November (Figures 5 and 6) and DJF (Figures 7 and 8). The reason for focusing on these two periods is to further examine the switch from weaker (in November) to stronger (in DJF) wave activity propagating into the lower stratosphere, as well as the strengthening and poleward-displacement of the polar front jet in the troposphere (in DJF).

3.3 A dynamical analysis of trends: November

Figure 5a shows the latitude-height profile of the zonally averaged eddy heat flux $v'T'$ climatology from ERA5 for November, which is dominated by negative values from 45-80°S in the lower stratosphere, consistent with upward propagating waves along the polar vortex edge. Quantitatively similar results can be obtained from the other three reanalyses (not shown). Figure 5 (c,e,g,i) shows the trend in eddy heat flux for November, which for all four reanalyses is marked by positive values in the lower stratosphere at 40-80°S, so in agreement with Figure 4 and confirming the reduction of poleward eddy heat flux / upward wave activity flux from the troposphere into the lower stratosphere. Overall, in terms of both magnitude and location, the best agreement is found between ERA5 and JRA-55, while the positive trend in CFSR is around half that of ERA5, indicating a much weaker reduction in upward wave activity from below for CFSR. This is despite CFSR showing stronger positive wind trends in the lower stratosphere compared to the other reanalyses in November (Figure 2), which is dynamically inconsistent as this would be expected to be associated with a relative stronger (rather than weaker) reduction in upward wave activity.

Figure 6 shows the 100-30 hPa vertically integrated trend (and climatology) of eddy heat flux for all waves, as well as planetary and synoptic waves, again for the month of November. This analysis confirms that the reduced upward wave fluxes in the lower stratosphere are composed of planetary waves, in good agreement with Orr et al. (2012). However, there is a large amount of disagreement in the CFSR trends compared to the other three reanalyses in terms of both amplitude and latitudinal extent.

Figure 5b shows the climatology of the November, zonally averaged eddy momentum flux $\overline{u'v'}$ derived from ERA5, which is dominated by negative values at 30-60°S in the lower stratosphere, indicating poleward momentum fluxes. In the troposphere, the climatology is marked by much larger negative values at 30-55°S and positive values at 60-80°S, indicating momentum

convergence in mid-latitudes. Figure 5 (d,f,h,j) shows the trend in eddy momentum flux, which for all four reanalyses at around 50-80°S is marked by negative values at ~100 hPa, so in agreement with Figure 4 and confirming enhanced poleward eddy momentum flux / equatorward propagation of wave activity. All four reanalyses show this feature, except that the magnitude of the trend is larger in MERRA-2 and even larger and more poleward in CFSR. Note that there are also positive trends at ~300 hPa, which are also apparent in Figure 4. Figure 6 (b,d,f) shows that the negative lower stratospheric trends displayed in Figure 5 are dominated by the contribution from planetary waves. Similar to the eddy heat fluxes, there is a large amount of disagreement in the CFSR trends compared to the other three reanalyses, while the best agreement is found between ERA5 and JRA-55.

3.4 A dynamical analysis of trends: Austral summer

Figure 7 is analogous to Figure 5, but for DJF. The eddy heat flux climatology for DJF from ERA5 (Figure 7a) is dominated by negative values at 40-60°S, 100-1000 hPa, indicating that upward propagating baroclinic waves are confined largely to the troposphere, as expected in austral summer (Plumb, 2011). Quantitatively similar results can be obtained from the other three reanalyses, with differences of only around $1 \text{ m s}^{-1} \text{ K}$ at a few locations (not shown). Figure 7 (c,e,g,i) shows results for the DJF trend for ERA5, JRA-55, MERRA-2, and CFSR. For all four reanalyses there is a negative trend poleward at around 50°S in the lower stratosphere (and to a lesser extent the uppermost region of the troposphere), so in agreement with Figure 4 and confirming the importance of enhanced upward wave fluxes at high latitudes into the lower stratosphere in the summer months (Orr et al., 2012). ERA5 and JRA-55 again show the best agreement, with MERRA-2 and especially CFSR showing larger negative values in the lower stratosphere (~300 hPa).

Figure 8 is analogous to Figure 6, but for DJF and the height range of 30-300 hPa for the eddy heat flux and 100-500 hPa for the eddy momentum flux. The reason for selecting different ranges for the vertical integration was because the strongest trends in eddy heat flux are found from 30-300 hPa for all four reanalyses, and from 100-500 hPa for the eddy momentum flux (Figure 7). Figure 8 (a,c,e) shows that the eddy heat flux trend from 30-300 hPa due to all waves is dominated by negative values at 45-80°S, which is poleward of the climatological values at 30-70°S (cf. Figure 7). In agreement with Orr et al. (2012), these trends are dominated by planetary waves at 55-80°S (Orr et al., 2012), while synoptic waves also have some role at 45-70°S. As the climatological tropopause height is above 300 hPa equatorward of 60°S (Figure 7(a,b)), some of the synoptic waves in this region are actually in the upper troposphere and not the lower stratosphere. Again, ERA5 and JRA-55 are in good agreement, while the MERRA-2 and CFSR trends are both stronger and more poleward.

Figure 7b shows the DJF eddy momentum flux climatology from ERA5. The climatology is marked by positive values at 60-75°S, 200-500 hPa and negative values at 30-55°S, 100-500 hPa, so confined largely to the troposphere. Similar climatologies can be obtained from the other three reanalyses with differences of no more than $4 \text{ m}^2 \text{ s}^{-2}$ at a few locations within the positive and negative regions shown for ERA5. Figure 7(d,f,h,j) shows DJF trends in momentum flux derived from ERA5, JRA-55,

MERRA-2, and CFSR. The trends are marked by negative values reaching $-5 \text{ m}^2 \text{ s}^{-2} \text{ dec}^{-1}$ in the troposphere at $40\text{-}70^\circ\text{S}$, so
330 consistent with Figure 4 and confirming the importance of enhanced poleward eddy momentum fluxes at the core of the
climatological polar front jet in the troposphere (Orr et al., 2012). All four reanalyses capture this feature, except that the
magnitude of the trend is largest in ERA5. The other three reanalyses produce the effect with a slightly more poleward shift.

Figure 8 (b,d,f) shows vertically integrated results for DJF from 100-500 hPa for the eddy momentum flux. The trends in eddy
momentum fluxes due to all waves are also dominated by negative values centered at $40\text{-}70^\circ\text{S}$ (cf. Figure 7), which is poleward
335 of the climatological minimum values and also dominated by the contribution from synoptic-scale waves. This is again in
agreement with Orr et al. (2012). The four reanalyses, however, exhibit more considerable disagreement in the trends that are
more pronounced than the differences in their climatological values.

3.5 Sensitivity of the trends to time period

To assess the statistical robustness of the trends, we explore the impact of small shifts in the time period of the analysis on the
trend. Figure 9 shows time-height cross-sections of the trends in zonally averaged zonal wind for the reanalyses from
340 September to February for three different 20-year periods (1980 to 1999, 1981 to 2000, and 1982 to 2001) that overlap our
analysis period of 1980 to 2001. The trends and spread in zonal wind between the reanalyses for the different periods agree
with the results for the 1980 to 2001 period. To examine the robustness of the trends in dynamical quantities, Figure 10
compares the spread of the November trends in 30 to 100 hPa vertically integrated eddy heat flux for the three 20-year periods
345 (c.f. Figure 6). The spread of the trends in eddy heat flux for the different periods are similar, and consistent with the results
for the 1980 to 2001 period. Examination of the sensitivity of the trends for the other dynamical quantities examined in this
study to the different time periods exhibited a similar robustness (not shown). The differences among the reanalyses are of
similar magnitude compared to the sampling uncertainty associated with the choice of time period. The choice of end points
does not seem to induce a systematic bias, e.g. towards smaller or larger trends in any of the reanalyses, or in the difference
350 between the reanalyses.

4 Discussion and summary

Differences in the formulation of reanalysis systems and their observational inputs can lead to significant differences in their
representation of the atmosphere, particularly for variables that are not directly observed (Fujiwara et al., 2017). Given the
relatively limited observations over Antarctica, there is greater potential for spread in their representation of the SH circulation
355 response to the ozone hole. Our results suggest that that there is nonetheless a high degree of consistency across the four
reanalysis datasets in the representation of the dynamical changes associated with ozone depletion. This conclusion is based
on a thorough assessment of trends in the zonally averaged zonal wind, eddy heat flux, eddy momentum flux, and wave forcing
(EPFD).

The expected strengthening of the lower stratospheric polar vortex during the austral spring-summer season and poleward shift of the polar front jet in the troposphere during summer is apparent in all four reanalyses. The differences in the trends in zonal wind between ERA5, JRA-55 and MERRA-2 is generally small in both the lower stratosphere and troposphere, with the largest differences of the order $0.2 \text{ m s}^{-1} \text{ dec}^{-1}$, which is small compared to the size of the reanalysis ensemble mean trends (up to $5 \text{ m s}^{-1} \text{ dec}^{-1}$ in the stratosphere and $2 \text{ m s}^{-1} \text{ dec}^{-1}$ in the troposphere). CSFR, however, shows greater disagreement compared to the other three reanalyses, evident by a relatively stronger wind increase in the lower stratosphere in spring and a larger poleward displacement of the polar front jet in summer (resulting in differences in the troposphere of up to $1 \text{ m s}^{-1} \text{ dec}^{-1}$).

The good agreement between ERA5 and JRA-55 circulation trends is perhaps because they both employ a 4D-VAR assimilation scheme, which is more sophisticated than the 3D-FGAT scheme employed by MERRA-2 and the 3D-VAR scheme employed by CSFR. This perhaps also explains why CSFR disagrees with the other three reanalyses in terms of the temperature trends, evident by an earlier onset to the cooling in the lower stratosphere in spring and enhanced cooling between 300 and 100 hPa throughout September to February. However, such disagreements could also depend on the observations that they assimilate (Manney et al., 2005; Lawrence et al., 2015). Long et al. (2017) shows that disagreements between reanalyses in the lower stratosphere temperature at SH high-latitudes are greater during the period 1979 to 1998 (corresponding to the assimilation of TIROS Operational Vertical Sounder (TOVS) data), which largely corresponds to the period examined in this study, and less afterwards during the ATOVS (Advanced TOVS) period from 1999 to 2014. The ability of each reanalysis to transition seamlessly between different satellite and other data sources at different times is also an issue, with more recent reanalysis having fewer discontinuities, while CSFR is characterized (globally) by multiple discontinuities in the stratosphere over the 1980 to 2001 period (Long et al., 2017). How the reanalysis systems include ozone and treat its radiative feedback also varies widely between reanalysis and might be an additional factor (Davies et al., 2017).

The circulation changes are consistent with our dynamical understanding of the stratosphere-troposphere system and are explainable in terms of four stages, which are apparent in all four reanalyses. An initial strengthening of the circulation in response to radiative cooling during the ‘onset’ stage plays an important role in conditioning the polar vortex so that fewer planetary waves can propagate up from the troposphere. The strengthening of stratospheric vortex winds in spring (mainly November) during the ‘growth’ stage is associated with a positive trend in EPFD. This coincides with reduced upward planetary wave activity fluxes at high latitudes from the troposphere into the lower stratosphere, causing a decrease in the wave-driven deceleration of the polar vortex. The weakening of the strengthened vortex in summer during the ‘decline’ and ‘decay’ stages is associated with a negative trend in EPFD. This coincides with increased upward planetary wave activity fluxes from the troposphere into the lower stratosphere at high latitudes due to the delay in the breakup of the stratospheric vortex, causing an increase in the wave-driven deceleration of the polar vortex. Both positive and negative trends in EPFD descend towards the tropopause, indicating a feedback between the strength of the vortex and the propagation of planetary waves (Chen and Robinson, 1992). The strengthening and poleward-displacement of the polar front jet in the troposphere

during the ‘decline’ and ‘decay’ stages are associated with changes to the synoptic-scale eddy fluxes of heat and momentum responsible for driving the tropospheric annular modes, which is evident by enhanced poleward eddy momentum fluxes into the jet. These changes in wave forcing and wave propagation are described by Orr et al. (2012, 2013), as well as other studies such as Hartmann et al. (2000), McLandress et al. (2010, 2011), and Hu et al. (2015). They agree with the temporal evolution of the zonal wind trends, although this does not necessarily indicate causality.

It is found that although the circulation trends are generally similar from one reanalysis to the next (with the exception of CSFR), significant discrepancies in the EPFD trends in the troposphere among the four reanalyses show up as alternating negative and positive horizontal stripes, which can be greater than the size of the mean trends across all reanalyses. Lu et al. (2014) suggest that the main contributor for such discrepancies are differences in the vertical derivative of the temperature, which are related to known issues with temperature increments caused by systematic biases in the assimilation of satellite measurements (e.g., Kobayashi et al., 2009). An additional factor could also be that derivatives are sensitive to interpolation from model levels to standard pressure levels. However, as there are no vertically alternating differences in the tropospheric wind trend, this suggests that this potential issue is relatively well constrained by analysis increments during data assimilation, while the EPFD is more model dependent. In the lower stratosphere, the trend in EPFD shows little difference among the four reanalyses.

The disparity between the size of the differences in wind trend and differences in eddy fluxes is also apparent. There are significant discrepancies in the associated trends in the eddy heat flux during the ‘growth’ stage (in November) and the ‘decline’ and ‘decay’ stages (in DJF) in the lower stratosphere, and the eddy momentum flux during the ‘decline’ and ‘decay’ stages in the troposphere. For CSFR, the positive trend in eddy heat flux during November is around half that of ERA5, indicating a much weaker reduction in upward wave activity / smaller reduction in wave-driven deceleration, despite it showing stronger positive wind trends in the lower stratosphere compared to the other reanalyses, which is dynamically inconsistent. This suggests that the eddy fluxes are also less constrained by the assimilation of observations, and that reanalysis temperature increments are able to cancel out differences in wave forcing, so that ultimately the impact on the large-scale circulation is small. Generally, across the four reanalyses, there is a large amount of disagreement in the CFSR wave forcing / propagation trends compared to the other three reanalyses, while the best agreement is found between ERA5 and JRA-55.

Another important source of possible dynamical inconsistency could stem from Coriolis torque on the residual meridional circulation and unresolved smaller scale forcing (Martineau et al., 2016), which although not considered in this study are both terms of the momentum budget (Eq. 1). Orr et al. (2012) investigated the role of the mean meridional circulation in the ozone hole momentum budget. They showed that the sum of the wave driving (EPFD) and Coriolis torque was in broad agreement with the actual zonal wind tendency. They further showed that the magnitude of the Coriolis torque was typically the same as the wave driving term, offsetting each other as expected under quasi-geostrophic scaling. Orr et al. (2012, 2013) also stress

that the circulation changes caused by the ozone hole are the result of both wave and radiative driving, although differences in radiative driving between the reanalyses are also not considered in this study.

425 To summarize, we show that all four modern reanalysis datasets provide a consistent estimate of the circulation changes due to the ozone hole, and that the discrepancies between the datasets are comparatively small. While our results show broad agreement on dynamical trends (eddy heat and momentum fluxes), there are non-trivial differences between reanalysis products, indicating that there is still room for improvement in our characterization of the atmosphere. Despite the consistency across reanalyses, it is possible that changes in the observational network over time could lead to spurious trends across them
430 all; they share the vast majority of the same input data. We have greater confidence in the trends in the circulation precisely because the changes can be explained by robust dynamical mechanisms. The reanalyses are both consistent with each other and self-consistent with our dynamical understanding of stratosphere-troposphere interactions. Looking forward, these findings will give us confidence that reanalysis datasets can be used to rigorously assess changes associated with the recovery of stratospheric ozone (Solomon et al., 2016; Banerjee et al., 2020), which is projected to return to 1980 levels within the next
435 few decades (Iglesias-Suarez et al., 2016).

Appendix A: Temperature trends

Temperature changes in the lower stratosphere are an important component of the ozone hole. To illustrate this, Figure A1 shows time-height cross-section of trends in zonally averaged temperature from September to February for ERA5, MERRA-2, JRA-55, and CFSR.

440 Code and data availability

The ERA5, JRA-55, MERRA-2, and CFSR zonal-mean data set of diagnostics used in this study are available for download from the CEDA (Centre for Environmental Data Analysis) website: <https://catalogue.ceda.ac.uk/uuid/b241a7f536a244749662360bd7839312> (Martineau, 2017).

Competing interests

445 The authors declare that they have no conflicts of interest.

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645 **Table 1. Key characteristics of ERA5, JRA-55, MERRA-2, and CSFR reanalysis systems. The abbreviations used are IFS (Integrated Forecast System), JMA GSM (Japanese Meteorological Agency Global Spectral Model), GEOS (Goddard Earth Observing System Model), NCEP CFS (National Center for Atmospheric Research Climate Forecast System), 3D-FGAT (Three Dimensional First Guess at Appropriate Time assimilation scheme), 4D-VAR (Four Dimensional Variational Data Assimilation), and 3D-VAR (Three Dimensional Variational Data Assimilation). In the column labelled ‘Model’ the year indicates the year for the version of the forecast model that was used for the reanalysis.**

Reanalysis	Reference	Model	Horizontal grid spacing	Assimilation scheme
ERA5	Hersbach et al. (2020)	IFS Cy41r2 (2016)	~ 31 km	4D- VAR (Hersbach et al., 2020)
JRA-55	Kobayashi et al. (2015)	JMA GSM (2009)	~ 55 km	4D- VAR (Kobayashi et al., 2015)
MERRA-2	Gelaro et al. (2017)	GEOS 5.12.4 (2015)	0.5°× 0.625°	3D-FGAT (Lawless et al., 2010)
CSFR	Saha et al. (2010, 2014)	NCEP CFS (2007)	0.3125°	3D- VAR (Saha et al., 2010, 2014)

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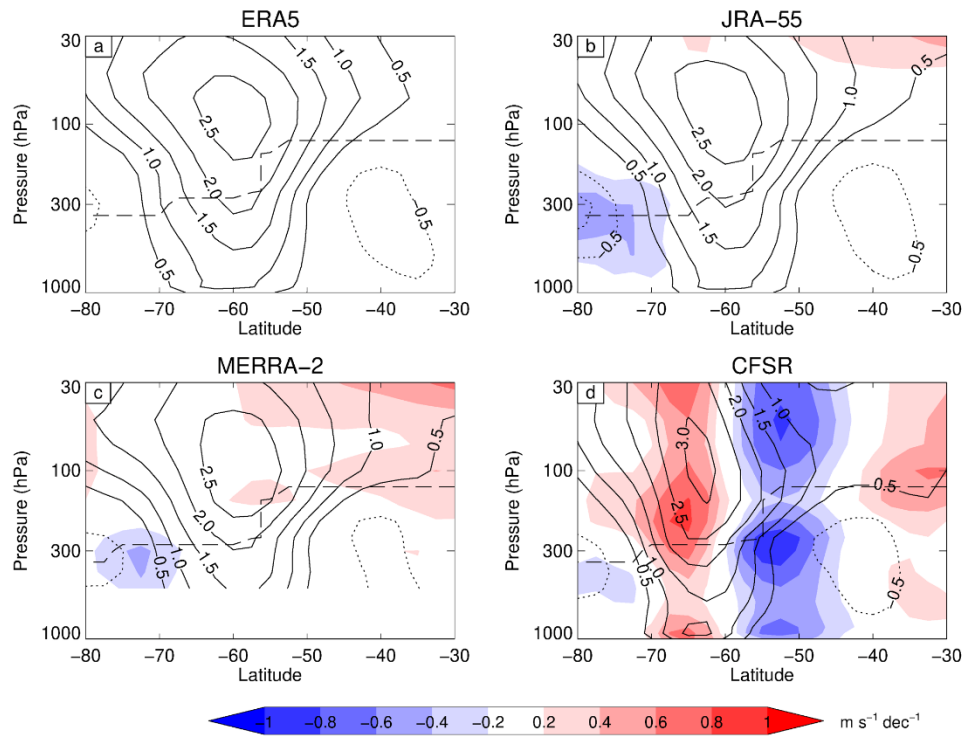


Figure 1: DJF trend of the zonally averaged zonal wind (contour intervals: ± 0.5 , ± 1.0 , ± 1.5 , ± 2.0 , ± 2.5 , ± 3.0 $\text{m s}^{-1} \text{dec}^{-1}$) from 1980 to 2001 for ERA5 (a), JRA-55 (b), MERRA-2 (c), and CFSR (d). The shading represents differences from ERA5 at intervals of ± 0.2 , ± 0.4 , ± 0.6 , ± 0.8 , ± 1.0 $\text{m s}^{-1} \text{dec}^{-1}$. The dashed line shows the climatological tropopause level. Note that results in the range 500 to 1000 hPa are not included in panel (c).

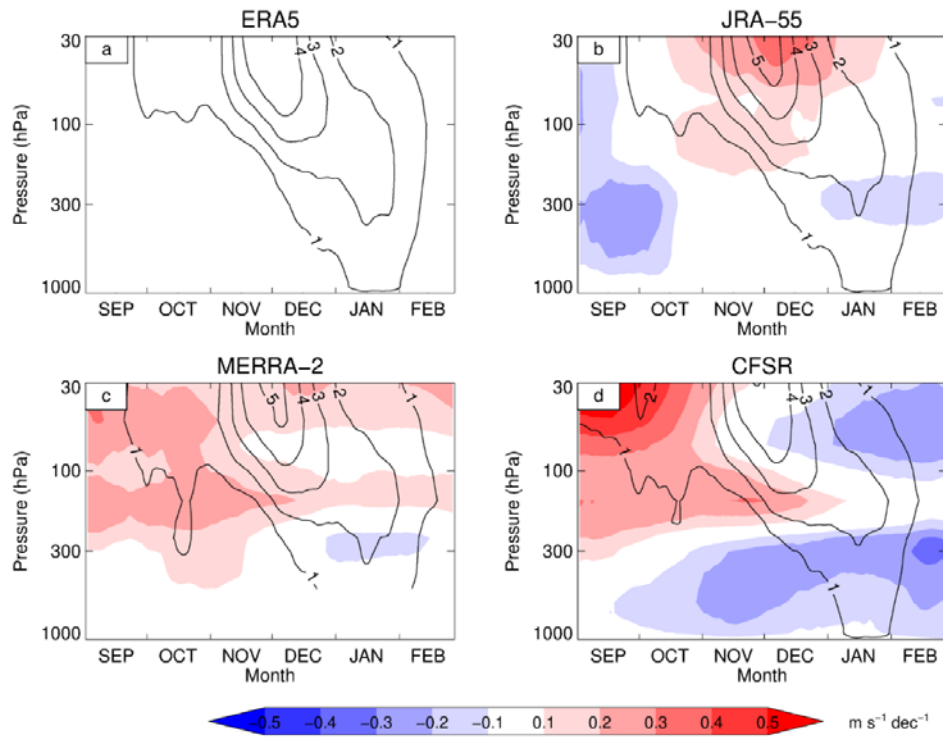
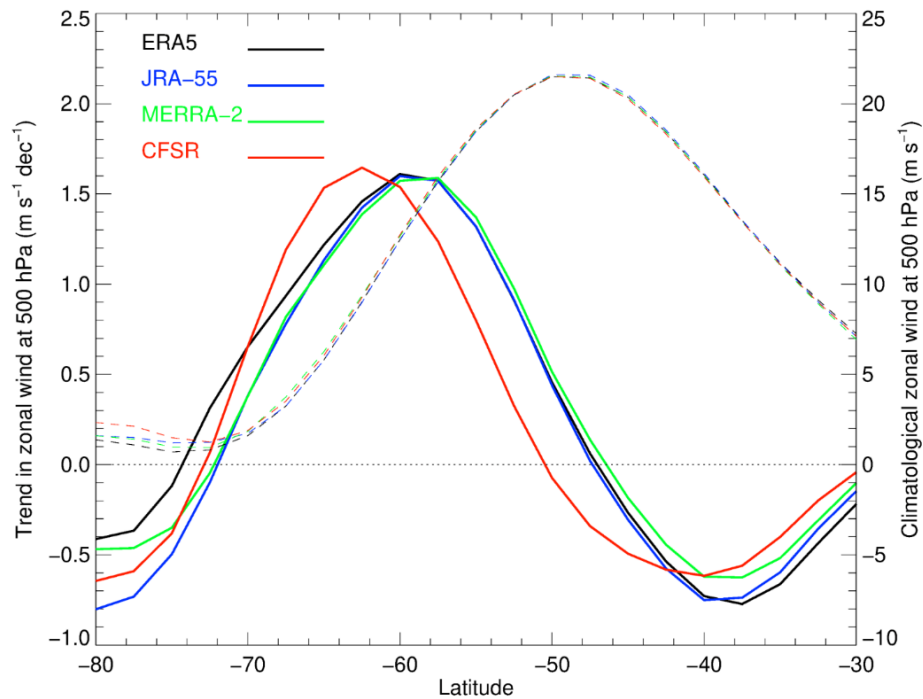


Figure 2: Time-height cross section of the trend in the zonally averaged zonal wind (contour intervals: 1, 2, 3, 4, 5 $\text{m s}^{-1} \text{dec}^{-1}$) averaged over 50 to 70°S from 1980 to 2001 for ERA5 (a), JRA-55 (b), MERRA-2 (c), and CFSR (d). The shading represents differences from ERA5 at intervals of ± 0.1 , ± 0.2 , ± 0.3 , ± 0.4 , ± 0.5 $\text{m s}^{-1} \text{dec}^{-1}$. Results in the range 500 to 1000 hPa are not included in panel (c). Note that for each panel the time-series is smoothed.



670 **Figure 3: DJF trend and mean in zonally averaged 500 hPa zonal wind from 1980 to 2001 for ERA5 (black line), JRA-55 (blue line),**
MERRA-2 (green line), and CFSR (red line). The trend is indicated by the thick lines (left y axis; units: m s⁻¹ dec⁻¹) and the
climatological mean by the thin dashed lines (right y axis; units: m s⁻¹). Note that the right and left axes have different scales.

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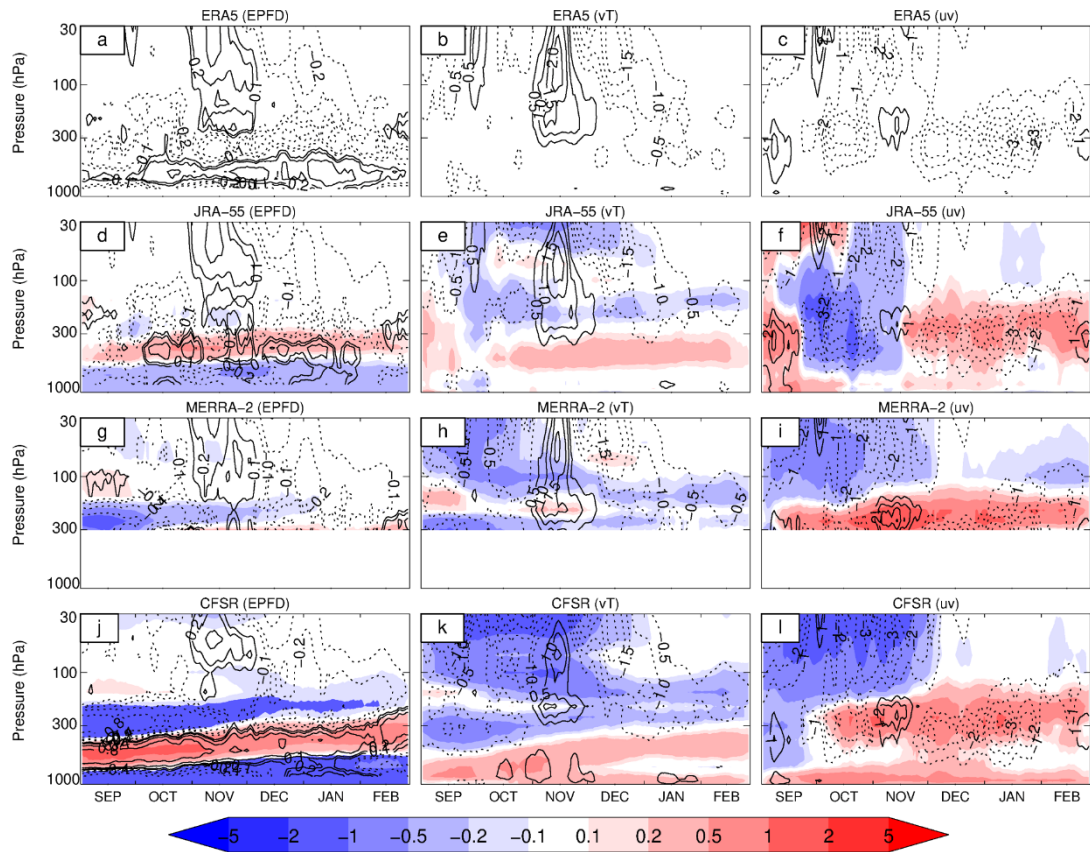


Figure 4: Time-height cross section of the trends in EP flux divergence (contour units: $\pm 0.1, \pm 0.2, \pm 0.4, \pm 0.8 \text{ m s}^{-1} \text{ d}^{-1} \text{ dec}^{-1}$; left column), eddy heat flux $\overline{v'T'}$ (contour units: $\pm 0.5, \pm 1.0, \pm 1.5, \pm 2.0, \pm 2.5, \pm 3.0 \text{ m s}^{-1} \text{ K dec}^{-1}$; middle column), and eddy momentum flux $\overline{u'v'}$ (contour units: $\pm 1, \pm 2, \pm 3, \pm 4, \pm 5 \text{ m}^2 \text{ s}^{-2} \text{ dec}^{-1}$; right column) averaged over 40 to 80°S from 1980 to 2001 for ERA5 (a, b, c), JRA-55 (d, e, f), MERRA-2 (g, h, i), and CFSR (j, k, l). The shading represents differences from ERA5 at intervals of $\pm 0.1, \pm 0.2, \pm 0.5, \pm 1.0, \pm 2.0, \pm 5.0$. Note that results in the range 300 to 1000 hPa are not included in panels (g), (h) and (i). Note that for each panel the time-series is smoothed.

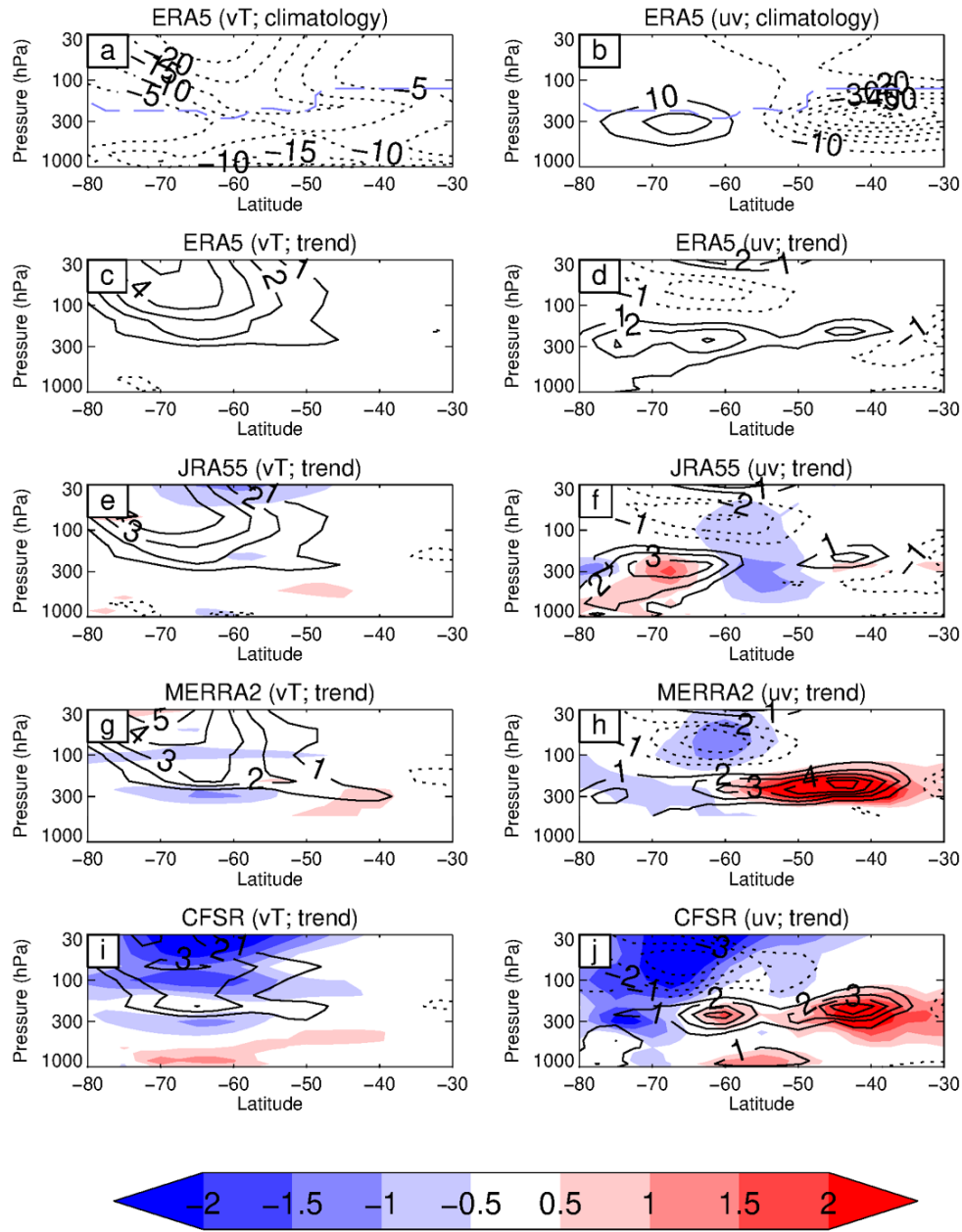


Figure 5: November trend of eddy heat flux $\overline{v'T'}$ (contour units: $\pm 1.0, \pm 2.0, \pm 3.0, \pm 4.0, \pm 5.0 \text{ m s}^{-1} \text{ K dec}^{-1}$; left column) and eddy momentum flux $\overline{u'v'}$ (contour units: $\pm 1.0, \pm 2.0, \pm 3.0, \pm 4.0 \text{ m}^2 \text{ s}^{-2} \text{ dec}^{-1}$; right column) from 1980 to 2001 for ERA5 (c, d), JRA-55 (e, f), MERRA-2 (g, h), and CFSR (i, j). The shading represents differences from ERA5 at intervals of $\pm 0.5, \pm 1.0, \pm 1.5, \pm 2.0$. Note that results in the range 500 to 1000 hPa are not included in panels (g, h). Panels (a, b) show the climatological mean values of $\overline{v'T'}$ ($\text{m s}^{-1} \text{ K}$) and $\overline{u'v'}$ ($\text{m}^2 \text{ s}^{-2}$) for ERA5 from 1980 to 2001, with the blue dashed line indicating the climatological tropopause level.

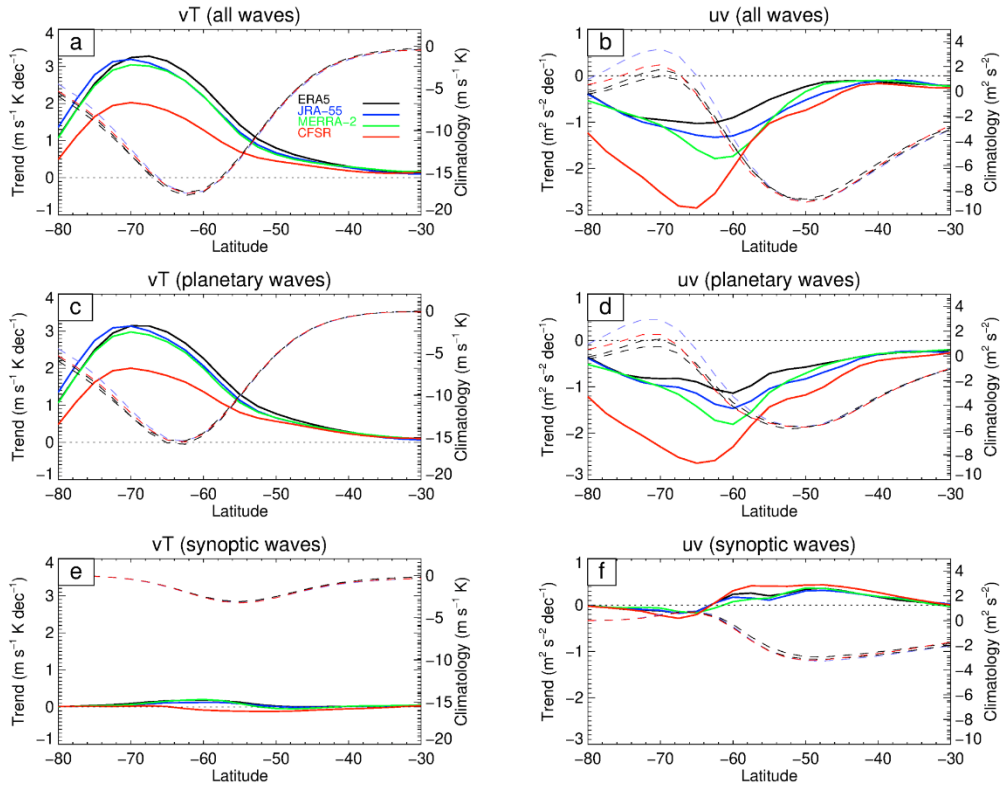
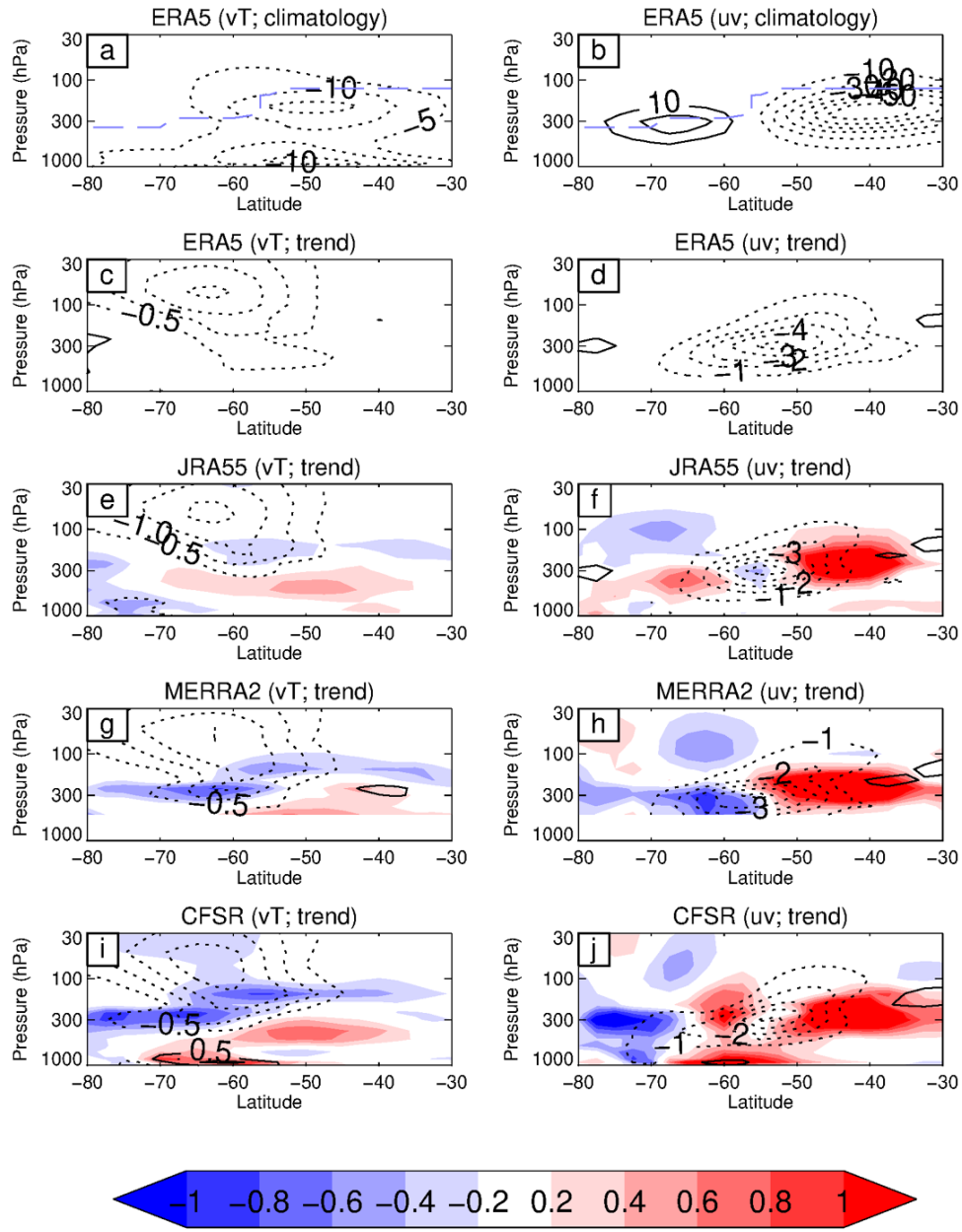
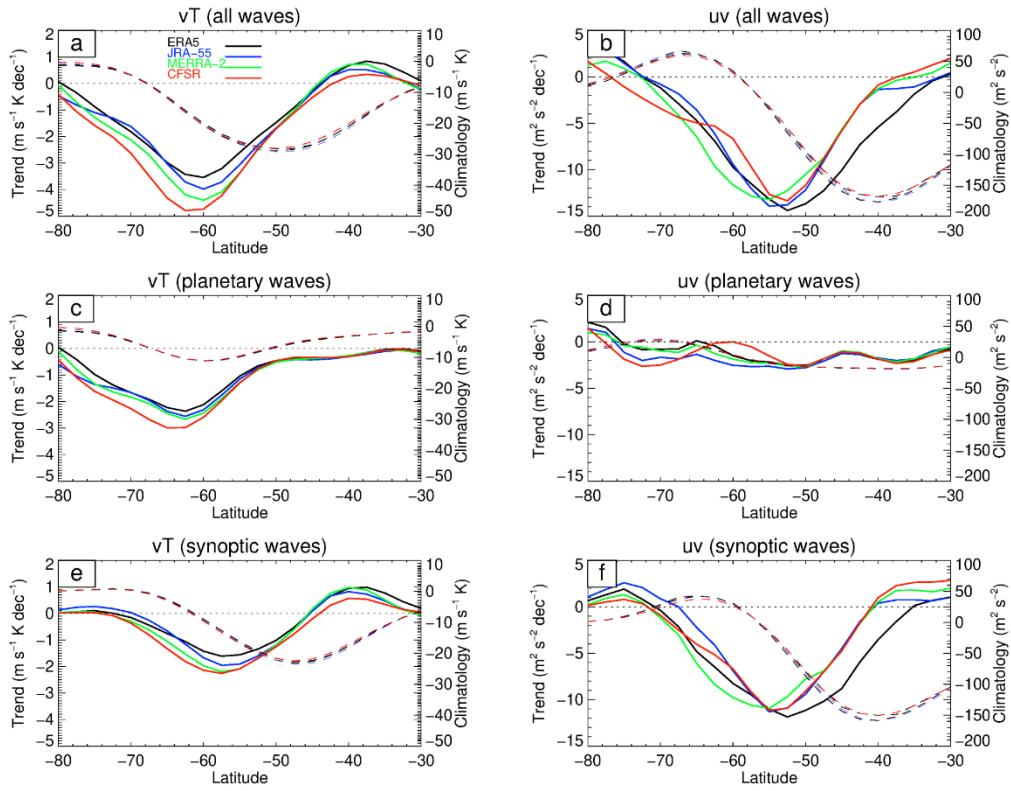


Figure 6: November trend and climatological mean in 30 to 100 hPa vertically integrated eddy heat flux $\overline{v'T'}$ (left column) and eddy momentum flux $\overline{u'v'}$ (right column) due to all waves (a, b), planetary waves (c, d) and synoptic waves (e, f) from 1980 to 2001 for ERA5 (black), JRA-55 (blue), MERRA-2 (green), and CFSR (red). The trend in $\overline{v'T'}$ is indicated by the thick lines (left y axis; units: $\text{m s}^{-1} \text{K dec}^{-1}$) and the climatological mean by the thin dashed lines (right y axis; units: $\text{m s}^{-1} \text{K}$). The trend in $\overline{u'v'}$ is indicated by the thick lines (left y axis; units: $\text{m}^2 \text{s}^{-2} \text{dec}^{-1}$) and the climatological mean by the thin dashed lines (right y axis; units: $\text{m}^2 \text{s}^{-2}$). Note that for both columns the right and left axes have different scales.



705 **Figure 7:** DJF trend of eddy heat flux $\overline{v'T'}$ (contour units: $\pm 0.5, \pm 1.0, \pm 1.5, \pm 2.0 \text{ m s}^{-1} \text{ K dec}^{-1}$; left column) and eddy momentum $\overline{u'v'}$ flux (contour units: $\pm 1.0, \pm 2.0, \pm 3.0, \pm 4.0, \pm 5.0, \pm 6.0 \text{ m}^2 \text{ s}^{-2} \text{ dec}^{-1}$; right column) from 1980 to 2001 for ERA5 (c, d), JRA-55 (e, f), MERRA-2 (g, h), and CFSR (i, j). The shading represents differences from ERA5 at intervals of $\pm 0.2, \pm 0.4, \pm 0.6, \pm 0.8, \pm 1.0$. Note that results in the range 500 to 1000 hPa are not included in panels (g, h). Panels (a, b) show the climatological mean values of $\overline{v'T'}$ ($\text{m s}^{-1} \text{ K}$) and $\overline{u'v'}$ ($\text{m}^2 \text{ s}^{-2}$) for ERA5 from 1980 to 2001, with the blue dashed line indicating the climatological tropopause level.



710 **Figure 8:** DJF trend and climatological mean in vertically integrated zonally eddy heat flux $\overline{v'T'}$ from 30 to 300 hPa (left column) and eddy momentum flux $\overline{u'v'}$ from 100 to 500 hPa (right column) due to all waves (a, b), planetary waves (c, d) and synoptic waves (e, f) from 1980 to 2001 for ERA5 (black), JRA-55 (blue), MERRA-2 (green), and CFSR (red). The trend in $\overline{v'T'}$ is indicated by the thick lines (left y axis; units: $\text{m s}^{-1} \text{K dec}^{-1}$) and the climatological mean by the thin dashed lines (right y axis; units: $\text{m s}^{-1} \text{K}$). The trend in $\overline{u'v'}$ is indicated by the thick lines (left y axis; units: $\text{m}^2 \text{s}^{-2} \text{dec}^{-1}$) and the climatological mean by the thin dashed lines (right y axis; units: $\text{m}^2 \text{s}^{-2}$). Note that for both columns the right and left axes have different scales.

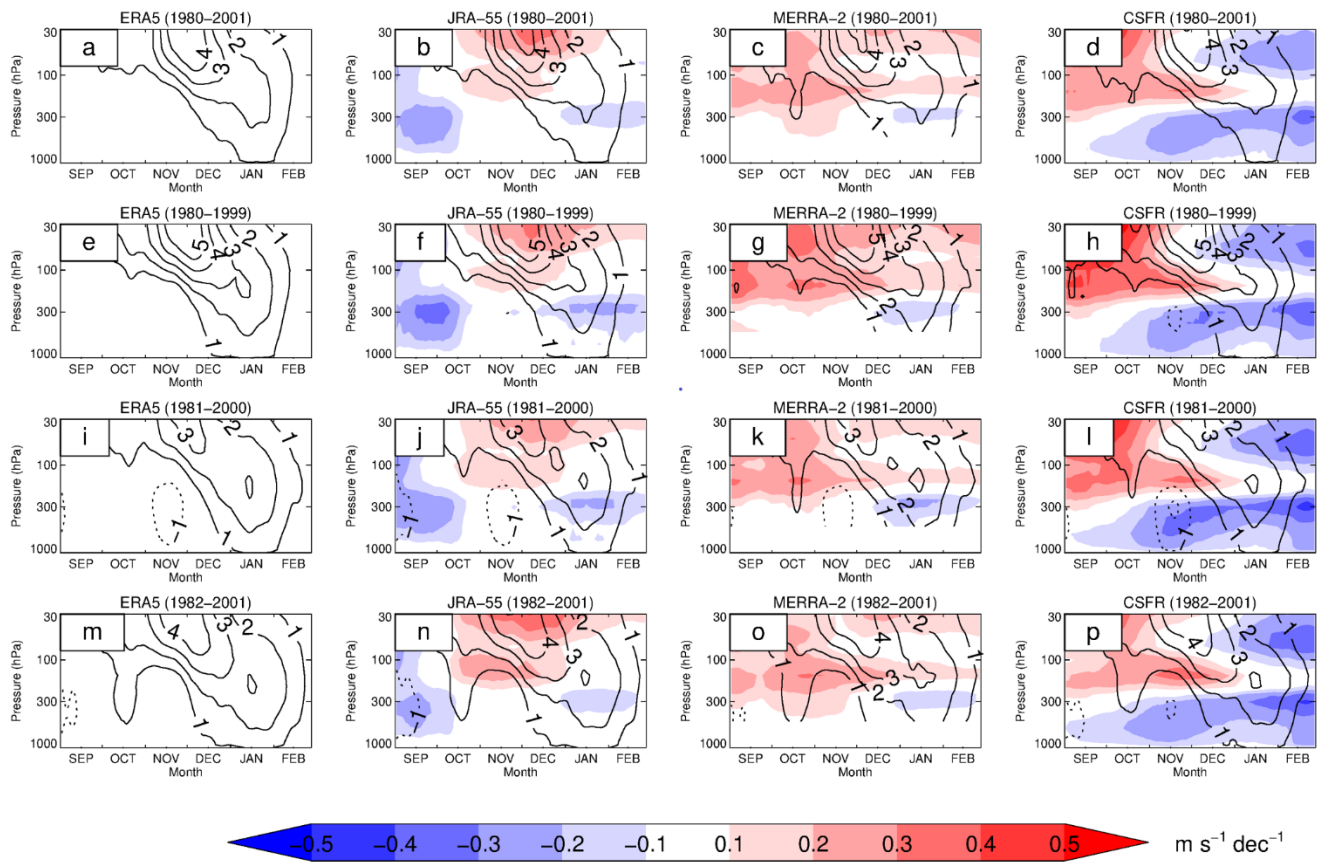
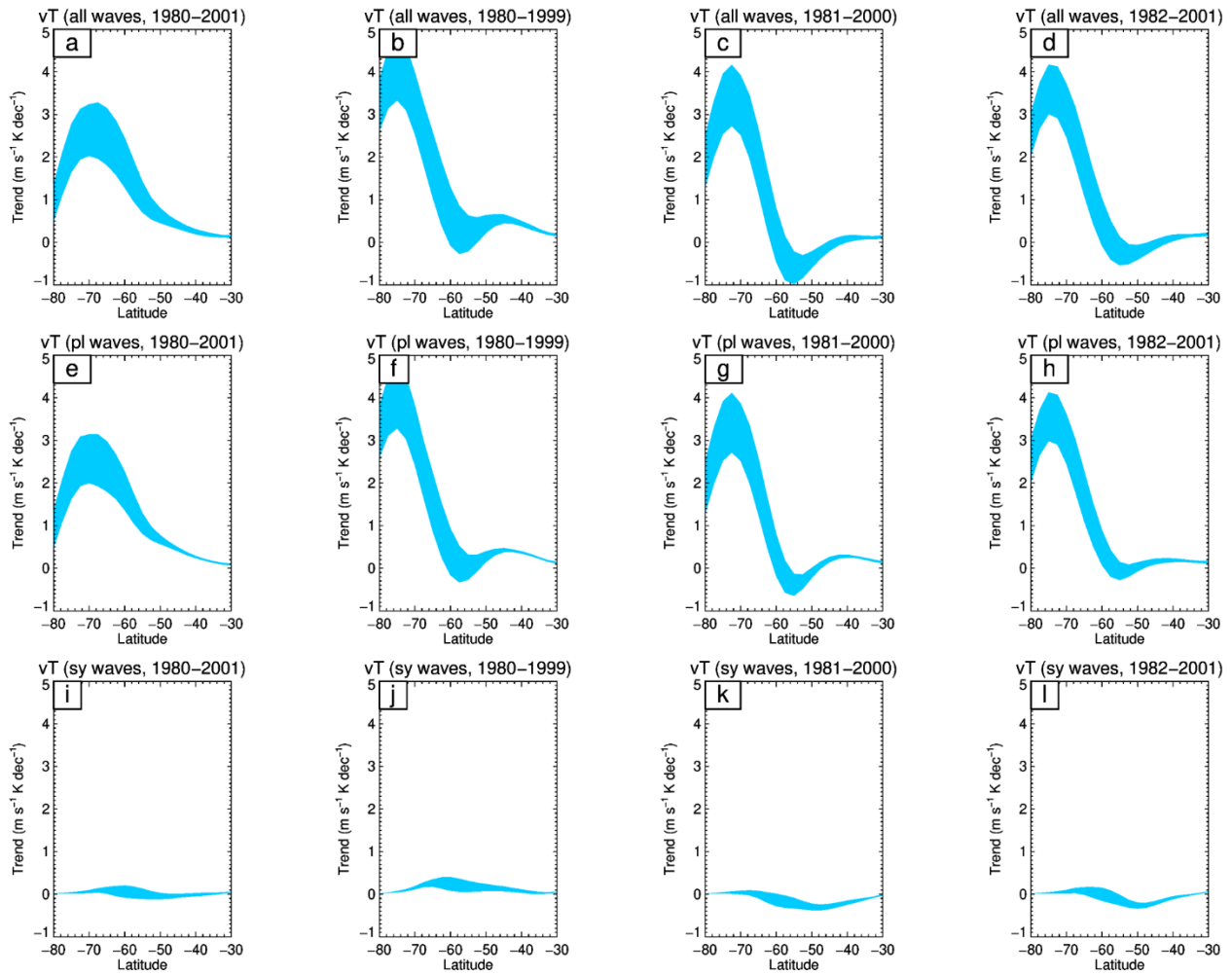


Figure 9: Sensitivity of the trend in zonal wind to time period, displayed as time-height cross section of the trend in the zonally averaged zonal wind (contour intervals: 1, 2, 3, 4, 5 $\text{m s}^{-1} \text{dec}^{-1}$) averaged over 50 to 70°S from 1980 to 2001 (a-d; same results as shown in Figure 2), 1980 to 1999 (e-h), 1981 to 2000 (i-l) and 1982 to 2001 (m-p) for ERA5, JRA-55, MERRA-2, and CSFR. The shading represents differences from ERA5 at intervals of ± 0.1 , ± 0.2 , ± 0.3 , ± 0.4 , $\pm 0.5 \text{ m s}^{-1} \text{dec}^{-1}$. Results in the range 500 to 1000 hPa are not included in panels (c, g, k, o). Panels (a-d) are the same results as shown in Figure 2. Note that for each panel the time-series is smoothed.



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Figure 10: Sensitivity of the trends in eddy heat flux to time period, displayed as shaded envelopes representing the spread (maximum and minimum values) derived from ERA5, JRA-55, MERRA-2 and CSFR of the November trend in 30 to 100 hPa vertically integrated eddy heat flux $\overline{v'T'}$ from 1980 to 2001 (a, e, i), 1980 to 1999 (b, f, j), 1981–2000 (c, g, k) and 1982–2001 (d, h, i) due to all waves (top row), planetary waves (middle row), and synoptic waves (bottom row). The units are $\text{m}^2 \text{s}^{-2} \text{dec}^{-1}$. Panels (a, e, i) are the same results as shown in the left column of Figure 6.

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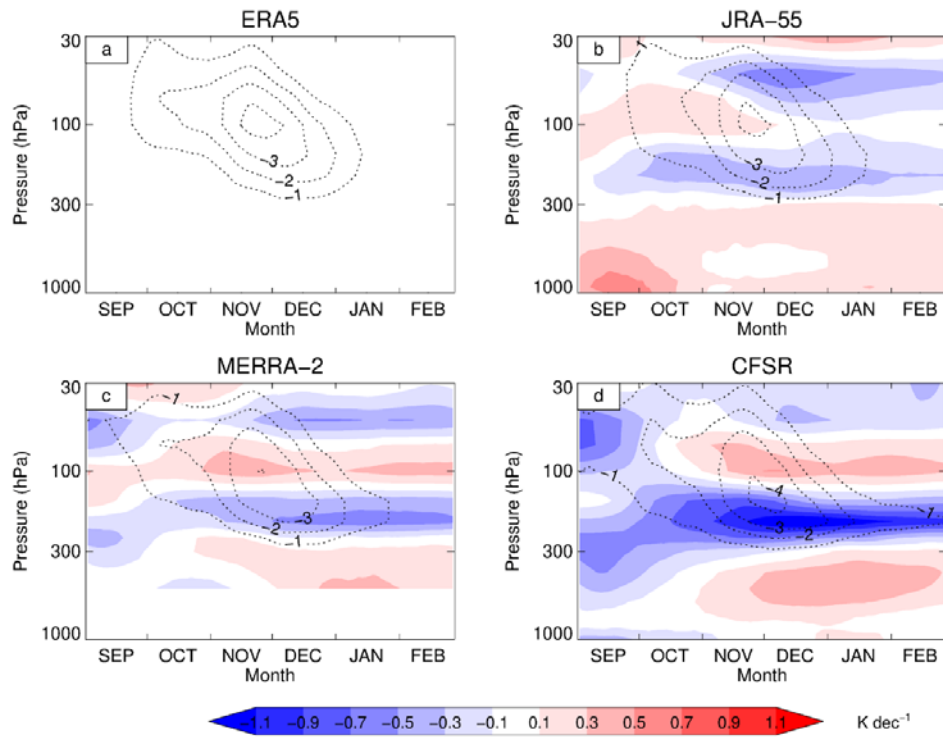


Figure A1: Time-height cross section of the trend in the zonally averaged temperature (contour intervals: -1, -2, -3, -4 K dec^{-1}) averaged over 70 to 87.5°S from 1980 to 2001 for ERA5 (a), JRA-55 (b), MERRA-2 (c), and CFSR (d). The shading represents differences from ERA5 at intervals of $\pm 0.1, \pm 0.3, \pm 0.5, \pm 0.7, \pm 0.9, \pm 1.1 \text{ m s}^{-1} \text{ dec}^{-1}$. Note that results in the range 1000 to 500 hPa are not included in panel (c). Note that for each panel the time-series is smoothed.