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Revealing the statistics of extreme events hidden in short weather forecast data

Justin Finkel¹, Edwin P. Gerber², Dorian S. Abbot³, Jonathan Weare²

¹Committee on Computational and Applied Mathematics, University of Chicago

²Courant Institute of Mathematical Sciences, New York University

³Department of Geophysical Sciences, University of Chicago

Key Points:

- Extreme weather risk, as measured by rate or return times, is inherently difficult to analyze because of data scarcity.
- Transition path theory reveals climatological statistics of sudden stratospheric warming events from high-fidelity subseasonal forecasts.
- Rates and seasonal distributions of 100-year stratospheric extremes are robustly computed from 47-day hindcast ensembles across 21 winters.

Corresponding author: Justin Finkel, justinfocus12@gmail.com

Abstract

Extreme weather events have significant consequences, dominating the impact of climate on society, but occur with small probabilities that are inherently difficult to compute. A rare event with a 100-year return period takes, on average, 100 years of simulation time to appear just once. Computational constraints limit the resolution of models used for such long integrations, but high resolution is necessary to resolve extreme event dynamics. We demonstrate a method to exploit short-term forecasts from a high-fidelity weather model and lasting only weeks rather than centuries, to estimate the long-term climatological statistics of rare events. Using only two decades of forecast data, we are able to robustly estimate return times on the centennial scale. We use the mathematical framework of transition path theory to compute the rate and seasonal distribution of sudden stratospheric warming (SSW) events of varying intensity. We find SSW rates consistent with those derived from reanalysis data, but with greater precision. Our method performs well even with simple feature spaces of moderate dimension, and holds potential for assessing extreme events beyond SSW, including heat waves and floods.

Plain Language Summary

Weather extremes are a continually recurring threat to human life, infrastructure, and economies. Yet, we only have sparse datasets of extremes, both simulated and observed, because by definition they occur rarely. We introduce an approach to extract reliable extreme event statistics from a non-traditional data source: short, high-resolution weather simulations. With 21 years of 47-day weather forecasts, we estimate probabilities of once-in-500-year events.

1 Introduction

The atmosphere’s extreme, irregular behavior is, in some ways, more important to characterize than its typical climatology. A society optimized for historical weather patterns is highly exposed to damage from extreme heat and cold, flooding, and other natural hazards. Extremes may respond more sensitively than mean behavior to climate change, an argument supported by elementary statistics (Wigley, 2009), empirical observations (Coumou & Rahmstorf, 2012; AghaKouchak et al., 2014; O’Gorman, 2012; Huntingford et al., 2014; Naveau et al., 2020) and simulations (Pfahl et al., 2017; Myhre et al., 2019). Recent unprecedented extreme weather events demonstrate the serious human impacts (Mishra & Shah, 2018; Van Oldenborgh et al., 2017; Goss et al., 2020; Fischer et al., 2021). The overall “climate sensitivity” (Hansen et al., 1984), summarized by a change in global-mean temperature, does not do justice to these consequences, which has led the community to develop “event-based storylines” (Shepherd et al., 2018; Sillmann et al., 2021) as a more tangible expression of climate risk.

The intermittency of extreme events makes precise risk assessment exceedingly difficult. 100 flips of a biased coin with $\mathbb{P}\{\text{Heads}\} = 0.01$ is almost as likely to yield zero heads (probability 0.366) as one head (probability 0.370), and half as likely to yield two heads (probability 0.185). Similarly, in a 100-year climate simulation or historical record, a once-per-century event may easily appear either non-existent or twice as likely as it really is. The difficulty exists even in a stationary climate, but worsens in the presence of time-dependent forcing, anthropogenic or otherwise. The limited historical record forces us to use numerical models as approximations, introducing a dilemma: we can run cheap, coarse-resolution models for long integrations, providing reliable statistics of a biased system, or expensive, high-resolution models for short integrations, which have lower bias but higher-variance due to under-sampling. Long-term climate simulations are usually performed with a low resolution of $O(50 - 100)$ km per grid cell (Haarsma et al., 2016). A coarse model might suffice to estimate global-mean temperature and other aggregated statistics, but cannot resolve convective systems, e.g., tropical cyclones and precipitation over complex topography, that deliver localized but heavy damage (O’Brien et al., 2016; He et al., 2019). Even large-scale events, such as a sudden stratospheric warming (SSW, the specific application of this paper) might arise from multi-scale interactions that are poorly represented in coarse model grids.

To obtain accurate dynamics and statistics, we must use the highest-fidelity models available, currently exemplified by the Integrated Forecast System (IFS) of the European Center for Medium-Range Weather Forecasts (ECMWF). Running at high resolutions of $\sim 16\text{-}32$ km (ECMWF, 2016), the IFS produces skillful ensemble forecasts spanning ~ 1 week-1 month. Such a high-resolution model can generate a highly plausible “storyline”, but cannot feasibly run long enough to estimate the climatological rate of an extreme event.

In this work, we help close this gap by assembling fragmented weather forecast ensembles together to cover the full dynamically relevant phase space. By re-weighting ensemble members in a principled way, we estimate probabilities of sudden stratospheric warming (SSW) events, in which the winter stratospheric polar vortex rapidly breaks down from its typical state, a strong cyclonic circulation over the winter-hemisphere pole. The associated subsidence and adiabatic warming can cause lower-stratospheric temperatures to rise by more than 40 K over several days (Baldwin et al., 2021). The reversal of stratospheric winds forces upward-propagating planetary waves to break at lower and lower levels, exerting a “downward influence” on tropospheric circulation (Baldwin & Dunkerton, 2001; Baldwin et al., 2003; Hitchcock & Simpson, 2014; Kidston et al., 2015). The midlatitude jet and storm track shift equatorward, bringing extreme cold spells and other anomalous weather to nearby regions (Kolstad et al., 2010; Kretschmer, Cohen, et al., 2018). King et al. (2019) documents the impact of an SSW on extreme winter weather over the British Isles, the so-called “Beast from the East” in February 2018. SSWs are a demonstrated source of surface weather predictability on the subseasonal-to-seasonal (S2S) timescale, a frontier of weather forecasting with many implications for helping humanity deal with meteorological extremes (Sigmund et al., 2013; Scaife et al., 2016; White et al., 2017; Vitart & Robertson, 2018; A. Butler et al., 2019; Lang et al., 2020; Bloomfield et al., 2021; Scaife et al., 2022). For these reasons, there is keen interest in improving (i) the prediction of SSW itself beyond the horizon of ~ 10 days that marks the current state-of-the-art (Tripathi et al., 2016; Domeisen et al., 2020), and (ii) understanding of the long-term frequency, seasonal distribution, and other climatological statistics of SSW.

The ensemble forecasts archived in the S2S project at ECMWF (Vitart et al., 2017) have the potential to provide more precise statistics than the limited historical data. We describe our data sources in section 2. To realize this potential requires a method to stitch the short trajectories together, which we outline in section 3 and describe more fully in Supporting Information. Section 4 presents our main result: with data consisting of 47-day forecasts over a 21-year period, we estimate rates and seasonal distributions of SSW events which, depending on severity, occur as rarely as once in 500 years. We discuss the implications in section 5 and conclude in section 6.

2 Data and definitions

Fig. 1(a,b) show the evolution of zonal-mean zonal wind at 10 hPa and 60°N (which we abbreviate $U_{10,60}$), a standard index for the strength of the stratospheric polar vortex. Black time-series show $U_{10,60}$ through two consecutive winters where SSW occurred, 2008-2009 (a) and 2009-2010 (b), superimposed on its 70-year climatology in gray from the ERA-5 reanalysis dataset (Hersbach et al., 2020). $U_{10,60}$ is typically positive throughout the winter months, characterizing a strong circumpolar jet that forms in the stratosphere during the polar night. Occasionally, however, the vortex breaks down and $U_{10,60}$ reverses direction, becoming negative in the middle of winter. This is the standard definition of an SSW event (e.g., A. H. Butler et al., 2015), but it does not capture the range of intensities between events. Clearly, January 2009 achieved a much more negative $U_{10,60}$ level than February 2010. More intense SSW events have been linked to stronger tropospheric impacts (Karpechko et al., 2017; Baldwin et al., 2021), which motivates our efforts to distinguish between them. Historical data can provide reasonably robust estimates of moderately rare events such as February 2010, in which $U_{10,60}$ barely reversed sign; events of this magnitude occur on average every two years. On the other hand, extraordinary events like January 2009 are quite poorly constrained due to small sample size, while carrying an outsize risk in a non-stationary climate (Fischer et al., 2021).

Figure 1. Climatology of polar vortex and illustration of dataset. (a,b): 70-year climatology of $U_{10,60}$ according to ERA-5, with the middle 40-, 80-, and 100-percentile envelopes in lightening gray envelopes. Two individual years are shown in black: 2008-2009 (a) and 2009-2010 (b). Two ensembles of S2S hindcasts (purple) are shown each winter, a small sample from the large S2S dataset of two ensembles *per week* from the ECMWF IFS. A range of SSW thresholds $U_{10,60}^{(th)}$ from 0 m/s to -35 m/s are marked by horizontal red lines. When $U_{10,60}$ crosses this line from above, an SSW has occurred, provided it happens between the vertical blue lines marking November 1 and Feb. 28. (c) Schematic of the Markov state model approximation we use to estimate rates. Blue and orange curves represent the partial trajectories from S2S. At each time step the data are clustered into discrete boxes, and probability transition matrices estimated by counting transitions from one day to the next.

116 To quantify SSW intensity, we vary the the $U_{10,60}$ threshold—henceforth called $U_{10,60}^{(th)}$ —
 117 from 0 m/s to -35 m/s in 5 m/s increments and consider each case separately. Horan and Reich-
 118 ler (2017) and A. H. Butler and Gerber (2018) have suggested the utility of examining different
 119 thresholds, as SSW events form a continuum. Horizontal red lines in Fig. 1(a,b) mark each thresh-
 120 old. Vertical blue lines frame the winter period of November 1-February 28 in which we allow
 121 SSWs to occur, to exclude “final warmings” at winter’s end when the vortex dissipates for the
 122 summer (Black et al., 2006). We only count the first event of the season, to avoid counting the
 123 subsequent oscillations of $U_{10,60}$ about $U_{10,60}^{(th)}$ as separate SSW events. A minimum separation
 124 time can also be imposed, as in (Charlton & Polvani, 2007), to allow multiple SSWs in a season,
 125 but these are rare and for the purpose of demonstration, we keep the definition as simple as pos-
 126 sible.

127 In addition to reanalysis, panels (a,b) also display a small sample of the S2S dataset in pur-
 128 ple. These are not forecasts but *reforecasts*, or *hindcasts*, generated by initializing a present-day
 129 model version on past weather conditions. The S2S archive compiles forecasts and hindcasts from
 130 11 forecasting centers around the world (Vitart et al., 2017), with a principle goal of tracking im-
 131 provements in skill from one version to the next. In this study we restrict ourselves to the ECMWF
 132 IFS, although our methodology can be repeated on other S2S datasets for intercomparison. We
 133 use data from the 2017 model version CY43R1, which produced 21 full winters of hindcasts be-
 134 tween autumn of 1996 and spring of 2017. These are initialized using ERA-Interim (ERA-I) re-
 135 analysis (ECMWF, 2011), which is almost identical to the more advanced ERA-5 from the stand-
 136 point of $U_{10,60}$. Two ensembles are launched every week, each with eleven members (one con-
 137 trol and ten perturbed forecasts) that run for 47 days before terminating. We use only the ten per-
 138 turbed members, which are initialized using a singular vector method and integrated with stochas-
 139 tic physics schemes (ECMWF, 2016). This introduces randomness into the ensemble, causing
 140 the members to drift apart over time after the initialization date, as shown in Fig. 1(c,d) for two
 141 sample ensembles. The specific strategy for perturbation of initial conditions and stochastic physics
 142 is informed by chaotic dynamical systems theory and has been refined by decades of numerical
 143 experiments (Mureau et al., 1993; Rabier et al., 1996; Palmer et al., 1998; Gelaro et al., 1998;
 144 Leutbecher, 2005; Lawrence et al., 2009; Buizza et al., 1999; Palmer et al., 2009) aimed at re-
 145 ducing forecast error due to under-dispersion, especially in the face of oncoming flow regime tran-
 146 sitions (Trevisan et al., 2001). In total, the S2S dataset contains over 900 years of simulation time.
 147 Many of them reach farther into the negative- $U_{10,60}$ tails than reanalysis, allowing us to calcu-
 148 late otherwise inaccessible probabilities.

149 3 Long-timescale dynamics from short trajectories

The advantage of sheer data volume comes with two attendant disadvantages. First, not all trajectories are independently sampled: on the contrary, all members of an ensemble are initialized close to reanalysis, and take several days to separate. Thus, the effective sample size is smaller

than 900 years. Second, no individual ensemble can directly provide an SSW probability beyond the 47-day time horizon, which is well short of the 120 days between November 1 and February 28 when SSWs are allowed to happen. To make use of the “hanging” trajectory endpoints and infer what might have transpired were the simulation to continue, we construct a *Markov state model* (MSM) (Deuffhard et al., 1999; Pande et al., 2010; Chodera & Noé, 2014) which is sketched in Fig. 1c. At every time sample $t = 1$ day, 2 days, ..., we partition state space into a disjoint collection of bins $S_{t,1}, S_{t,2}, \dots, S_{t,M_t}$ and approximate the transition probability matrix for each time-step from t to $t + 1$,

$$P_{t,t+1}(i, j) = \mathbb{P}\{\mathbf{X}(t+1) \in S_{t+1,j} | \mathbf{X}(t) \in S_{t,i}\}, \quad (1)$$

by counting the transitions between corresponding boxes. The matrices are row-normalized, which corrects for the redundancy and non-independence of ensemble members. Here, $\mathbf{X}(t)$ represents the full state vector of the ECMWF model. This sequence of matrices is the key ingredient that enables all downstream calculations, and it merits a brief note about the approximations involved. In a low-dimensional space, the partition could be created with a regular grid. However, every snapshot from the IFS has millions of degrees of freedom, including temperature and wind velocity in (latitude, longitude, pressure)-regular voxels. Any attempt to represent the dynamics of all these variables using a model such as (1) would suffer from large statistical error. On the other hand, if we only attempt to represent the dynamics of a small set of variables, our approximations may be very biased. To balance these concerns, we build the sets $S_{t,i}$ using k -means clustering of our data on a feature space Φ consisting of time-delays of $U_{10,60}$:

$$\Phi(\mathbf{X}(t)) = [U_{10,60}(\mathbf{X}(t)), U_{10,60}(\mathbf{X}(t-1)), \dots, U_{10,60}(\mathbf{X}(t-\delta))] \quad (2)$$

150 where $\delta = 20$ days is the number of retained time-delays, which can range from 15 to 25 with
 151 only minor effects on the results. We have also experimented with richer feature spaces includ-
 152 ing EOFs of geopotential height, but found these unnecessary. A growing body of theoretical (Takens,
 153 1981; Kamb et al., 2020) and empirical (Broomhead & King, 1986; Giannakis & Majda, 2012;
 154 Brunton et al., 2017; Thiede et al., 2019; Strahan et al., 2021) evidence supports the use of time-
 155 delay coordinates as reliable features for related methods. The k -means clustering is carried out
 156 using `scikit-learn` (Pedregosa et al., 2011) with $k = M_t$ on the collection of hindcast trajec-
 157 tories that were running between days t and $t + 1$. The number of clusters is set to $M_t = 170$ or
 158 the number of data points available on day t , whichever is smaller.

159 We use *transition path theory* (TPT) as a framework for combining several key forecast
 160 functions (both forward and backward-in-time) to compute the steady-state statistics of rare tran-
 161 sition events (Vanden-Eijnden, 2014; Finkel et al., 2020; Miron et al., 2021; Finkel et al., 2021a).
 162 TPT is most often applied in molecular dynamics applications (Noé et al., 2009; Meng et al., 2016;
 163 Strahan et al., 2021; Antoszewski et al., 2021) and is typically formulated in a time-homogeneous
 164 setting. The different timescales of climate applications, in particular the seasonal cycle, demand
 165 incorporating time-dependence explicitly, which we do in a manner similar to (Helfmann et al.,
 166 2020). Supporting Information provides more detail on TPT. All of the key forecast functions
 167 can be estimated directly using the transition matrix described above. In fact, the forecast func-
 168 tions each solve an infinite dimensional Feynman-Kac equation involving the transition opera-
 169 tor of the process (Strahan et al., 2021), and our partitioning of space into clusters corresponds
 170 to a basis expansion approach to solving those equations. This more general perspective moti-
 171 vates the *dynamical Galerkin approximation* (DGA) method of which our MSM approach is a
 172 special case (Thiede et al., 2019; Strahan et al., 2021; Finkel et al., 2021b, 2021a). MSMs are sim-
 173 ilar in spirit to analogue forecasting (van den Dool, 1989), which is enjoying a renaissance with
 174 novel data-driven techniques, especially for characterizing extreme weather (Chattopadhyay et
 175 al., 2020; Lucente et al., 2021). Formally, the transition operator encoded by the matrix in (1)
 176 is related to linear inverse models (LIMs; Penland & Sardeshmukh, 1995), which have also been
 177 used to predict atmospheric rivers at the subseasonal timescale (Tseng et al., 2021). Both MSMs
 178 and LIMs are finite-dimensional approximations of the Koopman operator (Mezić, 2013; Mezić,
 179 2005; Klus et al., 2018). For TPT analysis, however, an MSM is more convenient, which is ex-
 180 plained in Supporting Information.

181 Detailed comparison in the following section reveals that the approach sketched here is sta-
 182 tistically consistent with the direct method of sample-averaging over historical SSW events from
 183 reanalysis. However, the MSM approach provides more precise estimates for the rarest of events
 184 like the SSW of January 2009.

185 4 Results

186 4.1 Rate estimates

187 Fig. 2 shows rate estimates computed from the S2S dataset using the MSM-based approach
 188 outlined in the previous section, as well as from several reanalysis datasets using the direct count-
 189 ing method. Each circle indicates a point estimate using all the data from a given source and times-
 190 pan. In the case of S2S (red) the circle shows the mean rate from five independent trials with dif-
 191 ferent seeds for k -means clustering. The thick and thin vertical lines represent the 50% and 90%
 192 confidence intervals respectively, estimated from the pivotal bootstrap procedure (Wasserman,
 193 2004). We treat a full winter as a single unit of data for resampling, and we resample 40 times
 194 with replacement to estimate error bars. Any error bar that reaches the bottom edge of the log-
 195 arithmic plot is understood to include zero.

196 Different reanalysis datasets have different strengths for comparison with S2S. The most
 197 direct comes from ERA-5 (1996-2016)—meaning winter 1996/7-winter 2016/7, inclusive, the
 198 same time period as the S2S data—shown in orange. The S2S integrations from CY43R1 were
 199 initialized from ERA-I rather than ERA-5, but $U_{10,60}$ is virtually identical in both products (see
 200 Fig. ??). ERA-5 (1996-2016) is an appropriate baseline to compare with S2S, as both make use
 201 of the same observations. The key difference is that our MSM makes use of all the S2S hindcast
 202 integrations as well. Across the range of $U_{10,60}^{(th)}$, the S2S rate is less than or equal to the ERA-5
 203 (1996-2016) rate. However, this does not mean the two results are statistically inconsistent: 21
 204 flips of a fair coin can yield a range of outcomes, with 6-8 heads (combined probability 0.18) oc-
 205 ccurring slightly more often than either of the two most-likely outcomes of 10 or 11 heads (prob-
 206 ability 0.17 each). The orange error bars in Fig. 2 show the 50% and 95% confidence intervals
 207 of $(K/21)$, where K is a binomial random variable with $n = 21$ and $p =$ (the corresponding S2S
 208 estimate). In other words, we treat the S2S estimate as a null hypothesis and consider the real world
 209 as a sequence of independent draws from a probability distribution. For $U_{10,60}^{(th)} = -15$ m/s and
 210 above, the 21-year ERA-5 (1996-2016) point estimates are well within the 50% S2S confidence
 211 intervals, i.e., the interquartile range of $K/21$. For the more extreme events, the two estimates re-
 212 main consistent with 95%-level statistical significance, but ERA-5 (1996-2016) systematically
 213 indicates a higher frequency of extreme events in this 21-year timespan.

214 What climatology, then, is our MSM rate estimate inferring? Strictly speaking, it is a mix-
 215 ture between (i) the portion of phase space covered by 1996-2016 observations, and (ii) the *model*
 216 *climatology implied by the IFS*, including its stochastic parameterizations. Several recent stud-
 217 ies have performed the same task of filling out a sparse climate distribution using models (Horan
 218 & Reichler, 2017; Kelder et al., 2020), but with uninterrupted long runs of a global climate model.
 219 Our technique is novel in using short runs of a weather model instead.

220 Does the IFS climatology then correspond to anything in the real world? We can answer
 221 this by comparing to longer reanalyses, such as the 70-year ERA-5 (1950-2019) shown in gray
 222 in Fig. 2. Results are encouraging: ERA-5 (1950-2019) agrees with S2S in estimating a rate sys-
 223 tematically lower than ERA-5 (1996-2016), in other words suggesting this was an historically
 224 anomalously cold period. This tentative trend has been documented, and may explain some increasing
 225 cold-weather outbreaks despite an overall warming planet (Kretschmer, Coumou, et al., 2018;
 226 Garfinkel et al., 2017). Some studies indicate multi-decadal-scale variations in SSW frequency
 227 due to the quasi-biennial oscillation (QBO), El Niño southern oscillation (ENSO), Atlantic merid-
 228 ional overturning circulation, and other features of the coupled atmosphere-ocean system (Reichler
 229 et al., 2012; Dimdore-Miles et al., 2021). Hence, the recent barrage of SSWs may represent a tem-
 230 porary internal fluctuation rather than a secular trend. The consistency of S2S with ERA-5 on more

Figure 2. Rate estimates derived from S2S and reanalysis. Circles show point estimates of SSW rate according to each data source. S2S error bars show the 50% and 95% confidence intervals in thick and thin lines respectively, based on 40 bootstrap resamplings. Reanalysis error bars show the middle 50- and 95-percentile envelope of K/n , where K is a binomial random variable with p given by the corresponding S2S estimate, and n is the number of years in the reanalysis dataset. When an error bar overlaps with a reanalysis rate, the S2S rate is statistically consistent at the 95% confidence level.

Figure 3. Probability currents. The probability currents \mathbf{J}_{AB} (tendency of pre-SSW evolution) and \mathbf{J}_{AA} (tendency of non-SSW evolution) overlaid on the corresponding time-dependent probability densities π_{AB} and π_{AA} . Horizontal dashed line shows the boundary of B . The flux density of \mathbf{J}_{AB} across ∂B gives the seasonal distribution shown in Fig. 4.

231 common events, and the improvement of consistency with record length, is an encouraging sig-
 232 nal that the MSM estimate is extracting a meaningful statistic from the S2S dataset. This lends
 233 confidence in the S2S estimate as we reach farther into the negative $U_{10,60}$ tail where reanalysis
 234 data are too sparse to give any rate estimate.

235 Longer reanalysis is helpful to generate better statistics. For this, we incorporate one more
 236 relevant product, ERA-20C, which spans the longer period 1900-2007, but assimilates only sur-
 237 face measurements as opposed to satellite data (Poli et al., 2016). With these deliberate limita-
 238 tions, ERA-20C likely suffers higher bias than ERA-5 or ERA-I, but it enjoys lower variance due
 239 to its longer timespan. In their period of overlap (1950-2007, see Fig. ??), they roughly agree on
 240 the SSW rates with moderate thresholds of $U_{10,60}^{(th)} = 0$ and $U_{10,60}^{(th)} = -5$ m/s, but otherwise ERA-
 241 20C appears biased toward fewer SSW events. Nonetheless, ERA-20C is our best estimate for
 242 the SSW rate over the full 20th century.

243 In the upper range of thresholds from 0 m/s to -15 m/s, all datasets suggest a linear rela-
 244 tionship between $U_{10,60}^{(th)}$ and rate. In the lower range from -20 m/s to -35 m/s, reanalysis be-
 245 comes too noisy to discern clear trends, as these estimates rely on just a few exceptional events
 246 like January 2009 (Fig. 1). However, S2S clearly suggests an exponential trend with an e -folding
 247 scale of ~ 4 m/s. Events become tenfold rarer as the threshold is lowered by 10 m/s. These re-
 248 sults depend somewhat on parameter choices (see Supporting Information), but are robust to vari-
 249 ations in the delay time δ from 15 to 25 days.

250 4.2 Probability current

To explain the rate calculation, we briefly expand on the TPT framework, whose real strength is to not only provide numerical rates, but to decompose them into a sum over possible pathways into the rare event. The spread of pathways is encoded by the *probability current*, a vector field $\mathbf{J}_{AB}(t, \mathbf{x})$ over state space that indicates the average tendency of the system $\mathbf{X}(t)$ as it passes through state \mathbf{x} , conditioned on an SSW occurring. The subscript AB refers to two distinguished sets A and B in space-time,

$$A = \{(t, \mathbf{x}) : t < \text{Nov. 1 or } t > \text{Feb. 28}\} \quad (3)$$

$$B = \{(t, \mathbf{x}) : \text{Nov. 1} \leq t \leq \text{Feb. 28, and } U_{10,60}(\mathbf{x}) < U_{10,60}^{(th)}\}. \quad (4)$$

251 An SSW event can now be defined adhering to the TPT formalism (Vanden-Eijnden, 2014) as
 252 a passage of $\mathbf{X}(t)$ from A (the pre-winter part) to B , *before* returning to A (the post-winter part).
 253 Just as the symbol AB encodes an SSW, the symbol AA encodes a winter without SSW, in which
 254 the system departs A in the fall and re-enters A in the spring without ever hitting B . A second vec-
 255 tor field, $\mathbf{J}_{AA}(t, \mathbf{x})$, indicates the average tendency of the system during non-SSW winters. Both

256 currents, \mathbf{J}_{AB} and \mathbf{J}_{AA} , are computable in discretized forms from the transition matrices $P_{t,t+1}(i, j)$
 257 following Metzner et al. (2009). Consistent projections of these reactive currents from the full
 258 delay-embedded space down to U_{10} can be defined following Strahan et al. (2021) and are shown
 259 in Fig. 3. Supporting Information details the visualization procedure. The streamlines of \mathbf{J}_{AB} lead
 260 directly to the boundary ∂B of B , whereas the streamlines of \mathbf{J}_{AA} avoid this boundary and lead
 261 instead to ∂A (the right edge of the plot). Background shading indicates the corresponding time-
 262 dependent probability densities $\pi_{AB}(t, \mathbf{x})$ (a) and $\pi_{AA}(t, \mathbf{x})$ (b), defined as the density of all sys-
 263 tem trajectories $\mathbf{X}(t)$ destined for an SSW event or a non-SSW winter, respectively. Two sam-
 264 ples from each ensemble are superimposed: 1962-1963 and 2005-2006 as representative SSW
 265 winters, and 1966-1967 and 2004-2005 as representative non-SSW winters. The SSW trajecto-
 266 ries drop out of the ensemble when they first enter B and the curves turn from solid to dashed;
 267 this is why the total probability $\int \pi_{AB}(t, \mathbf{x}) d\mathbf{x}$ becomes steadily smaller as time progresses, be-
 268 cause it is an average over fewer and fewer events. In fact, one can show (see Supporting Infor-
 269 mation) that the $\pi_{AB}(t, \mathbf{x})$ is identical to the t -component of $\mathbf{J}_{AB}(t, \mathbf{x})$, which roughly quantifies
 270 how many SSW-bound trajectories are temporarily maintaining steady—or even increasing—
 271 $U_{10,60}$ before the upcoming event. Note that the individual trajectories do not track along stream-
 272 lines of the current: only their average evolution does. For example, the individual sample tra-
 273 jectories plummet toward B passing through *flat* \mathbf{J}_{AB} arrows, which account for the other SSW-
 274 bound trajectories that still persist at the same time of year.

These vector fields have concrete physical meaning: the field lines of \mathbf{J}_{AB} poke through ∂B
 with a time-dependent flux density that integrates to the total rate, as seen in the equation

$$\int_{\text{Nov. 1}}^{\text{Feb. 28}} \mathbf{J}_{AB} \cdot \mathbf{n} dt = \frac{\# \text{ SSW events}}{\text{Year}} \quad (5)$$

where \mathbf{n} is the unit vector in state space pointing directly into B ; in our case, $\mathbf{n} = -\nabla U_{10,60}(\mathbf{x}) / \|\nabla U_{10,60}(\mathbf{x})\|$.
 Moreover, SSW events can occur at different times during the winter, and the contribution from
 each time interval is equal to the corresponding partial flux integral. For example,

$$\int_{\text{Dec. 1}}^{\text{Dec. 31}} \mathbf{J}_{AB} \cdot \mathbf{n} dt = \frac{\# \text{ Dec. SSW events}}{\text{Year}} \quad (6)$$

275 This relation allows us to examine more refined details of SSW climatology: the seasonal dis-
 276 tribution of events.

277 4.3 Seasonal distribution

278 Past studies have found that seasonal differences are associated with dynamical differences
 279 in SSW events. For example, “Canadian warmings” shift the Aleutian high and occur earlier in
 280 the winter (A. H. Butler et al., 2015). Categorizing SSWs by their seasonality may reveal pre-
 281 ferred timings that indicate when and why the polar vortex is most vulnerable (Horan & Reich-
 282 ler, 2017). Unfortunately, month-by-month rate estimates from reanalysis are noisier than full-
 283 winter rate estimates, as splitting data into finer categories makes the events even sparser. We can
 284 again use S2S data to enhance precision by recruiting the larger database of partial trajectories.
 285 Fig. 4 shows seasonal distributions at two thresholds, $U_{10,60}^{(\text{th})} = -15$ m/s (left) and $U_{10,60}^{(\text{th})} = 0$ m/s
 286 (right), according to the same four datasets used in Fig. 2. Each panel displays the distribution
 287 at two resolutions: monthly (hashed) and sub-monthly (solid, and rounded to the nearest day),
 288 both according to the same dataset and with the same total integrals equal to the rate estimate.
 289 To express the seasonal cycle as a probability distribution, we normalize so that all histograms
 290 in Fig. 4 integrate to one, with units of probability per day. The two columns have different ver-
 291 tical scales to see features more readily.

292 Several features are noteworthy. For the conventional SSW, $U_{10,60}^{(\text{th})} = 0$ m/s, the reanaly-
 293 sis histograms all exhibit a common seasonal trend of steadily rising SSW frequency from Novem-
 294 ber to January and a small decline in February. The coarse S2S histogram disagrees, with a slight
 295 increase in February. Both trends are consistent with prior studies of seasonality at monthly res-
 296 olution (e.g., Charlton & Polvani, 2007). At a finer resolution of ~ 10 days, however, the S2S

Figure 4. Seasonal distributions of SSW events. Left and right columns show statistics with threshold $U_{10,60}^{(th)} = -15$ m/s and $U_{10,60}^{(th)} = 0$ m/s, respectively, and each row uses a different data source. Each panel has a hashed histogram at monthly resolution, along with a solid-colored histogram at $\frac{1}{3}$ -monthly resolution (rounded to the nearest day), with an equal area equal to unity. The vertical unit is SSW events per day. The vertical scales are shared within within each column, but different between columns in order to make the shape of the histogram at $U_{10,60}^{(th)} = -15$ m/s more easily visible.

297 histogram reveals a frequency peak in late January/early February and declines thereafter. The
 298 January/February peak is documented in the literature, e.g., by (Horan & Reichler, 2017), who
 299 diagnosed the peak as a balance between two time-varying signals: the background strength of
 300 the polar vortex, and the vertical flux of wave activity capable of disturbing the vortex. Addition-
 301 ally, the 10-day resolved S2S histogram reveals a smaller December peak, which is absent from
 302 ERA-5 reanalysis and at best noisily present in ERA-20C. The bimodal structure seen in S2S has
 303 also been found tentatively in prior studies with both reanalysis and models (e.g., Horan & Re-
 304 ichler, 2017; Ayarzagüena et al., 2019). We speculate that the early peak represents Canadian warm-
 305 ings (Meriwether & Gerrard, 2004), which our result suggests may deserve a more decisive clas-
 306 sification.

307 All three reanalysis-based estimates of SSW distributions have a low signal-to-noise ratio,
 308 exemplified by the intermittent frequency spikes. The hint of a third peak at the end of Febru-
 309 ary is clearer in reanalysis than S2S, and might be the beginning of the “final warmings”, but its
 310 significance is questionable because of the histograms’ general noisiness. This is even more of
 311 a problem at the more extreme threshold $U_{10,60}^{(th)} = -15$ m/s, where the ERA-5 (1996-2016) has
 312 degenerated to two isolated spikes while S2S retains a smoother shape, with little sign of bimodal-
 313 ity. Early December still supports a nonzero rate of extreme SSW events, but is not a highly fa-
 314 vorable time for them. This suggests that whatever distinct SSW type accounts for the Decem-
 315 ber peak at $U_{10,60}^{(th)} = 0$ m/s is limited to weaker events. These results are subject to all the caveats
 316 of our data-driven procedure (see Supporting Information), but merit further investigation with
 317 numerical models.

318 5 Discussion

319 By comparing S2S results with reanalysis, we are measuring the composition of three sep-
 320 arate error sources: (i) forecast model error, (ii) non-stationarity of the climate *with respect to SSW*
 321 *events* over the reanalysis period, and (iii) numerical errors in the MSM approach, both statis-
 322 tical (from the finite sample size) and systematic (from the projection of forecast functions onto
 323 a finite basis). We briefly address each error source in turn.

324 The S2S trajectories were realized only in simulation, not in the physical world. Accord-
 325 ingly, our S2S estimates apply strictly to the climatology of the 2017 IFS, a statistical ensemble
 326 that could be concretely realized by running the model uninterrupted for millennia, with exter-
 327 nal climatic parameters sampled from their variability in the short 21-year time window of 1996-
 328 2016. Such long, equilibrated simulations have been performed with coarser models by, e.g., Kelder
 329 et al. (2020) to assess UK flood risk (the so-called “UNSEEN” method), and by Horan and Re-
 330 ichler (2017) to assess SSW frequencies, but this is not practical given the constraints and mis-
 331 sion of the ECMWF IFS. Given these constraints, we have assembled our best approximation us-
 332 ing S2S trajectories. Indeed, the S2S dataset is an ensemble of opportunity for us. It was created
 333 to compare the skill of different forecast systems on S2S timescales, not at all for the purpose of
 334 establishing a climatology of SSWs.

335 The IFS model has proven outstanding in its medium-range forecast skill (Vitart, 2014; Kim
 336 et al., 2014; Vitart & Robertson, 2018). However, there is a caveat that the IFS was designed for

337 short forecasts, and it is not clear how it would behave if allowed to run for hundreds of years
338 as a climate model, which requires careful attention to the boundary condition and conservation
339 issues. Even if the climate were to remain stationary with its 1996-2016 parameters, numerical
340 and model errors would inject some bias into the equilibrated simulation. Repeatedly initializ-
341 ing S2S forecasts with reanalysis ensures a realistic background climatology, and allows us to
342 rely on the IFS strictly for the short-term integrations that it was designed for. Our method may
343 be used as a diagnostic tool to compare different models against each other, with specific atten-
344 tion paid to their rare event rates. A useful extension of this work would be to repeat the anal-
345 ysis on multiple data streams from all 11 forecasting centers worldwide that contribute to the S2S
346 project, providing a new rare event-oriented intercomparison metric.

347 The rate we estimate with an MSM is the SSW rate of the climate system frozen in its 1996-
348 2016 state. Comparing with a 70-year reanalysis dataset (ERA-5 1950-2019) measures the de-
349 parture of the 21-year SSW climatology from the 70-year climatology, and likewise for the 108-
350 year reanalysis ERA-20C (1900-2007). Of course, the 21-year SSW climatology itself may be es-
351 timated directly from reanalysis, but we have demonstrated in Fig. 2 that S2S gives more pre-
352 cise estimates that are different from the observations, but not at a statistically significant level.
353 Our results indicate that according to the 2017 IFS, 1996-2021 was more similar to 1950-2019
354 than direct counting of SSW events would suggest, which could of course mean that the IFS was
355 missing some key climatological variable during that period (Dimdore-Miles et al., 2021). There
356 is insufficient evidence on the anthropogenic influence on SSW to reject the hypothesis of sta-
357 tionarity (Ayarzagüena et al., 2020). By running our method on different historical periods, we
358 might discern a more decisive signal of secular changes than would be available from raw data.

359 Error source (iii) is the most open to scrutiny and improvement. In a sequence of preced-
360 ing papers (Finkel et al., 2021b, 2021a), we have benchmarked the performance of DGA (with
361 a similar MSM basis set) on a highly idealized SSW model due to Holton and Mass (1976). DGA
362 was originally developed in molecular dynamics to study protein folding and has been bench-
363 marked on a diverse set of low- and high-dimensional dynamical systems (Thiede et al., 2019;
364 Strahan et al., 2021; Antoszewski et al., 2021). Our parameter choices here, detailed further in
365 Supporting Information, are informed by prior experience. Nevertheless, large-scale atmospheric
366 models are a mostly-unexplored frontier for this class of methods. In this study, we have worked
367 with static datasets produced by some of the most advanced models in the world; however, an
368 even more powerful procedure would be to generate data adaptively.

369 Our method exceeds what is possible directly from reanalysis, but we are not yet fully “lib-
370 erated” from observations: every S2S trajectory is initialized near reanalysis, and it only has 47
371 days to explore state space before terminating. This fundamentally limits how far we can explore
372 the tail of the SSW distribution. In other words, the real climate system sets the “sampling mea-
373 sure” which is a flexible but important component in the DGA pipeline (Thiede et al., 2019; Stra-
374 han et al., 2021; Finkel et al., 2021b). On the other hand, with an executable model, we could ini-
375 tialize secondary and tertiary generations of short trajectories to push into more negative $U_{10,60}$
376 territory and maintain statistical power for increasingly extreme SSW events. This is the essence
377 of many rare-event sampling algorithms, such as those reviewed in Bouchet et al. (2019) and Sapsis
378 (2021). For example, a splitting large-deviation algorithm was used in Ragone et al. (2018) to
379 sample extreme European heat waves and estimate their return times. Quantile diffusion Monte
380 Carlo was used in Webber et al. (2019) to simulate intense hurricanes, and in (Abbot et al., 2021)
381 to estimate the probability of extreme orbital variations of Mercury. Many other rare event sam-
382 pling studies have been performed in fluid dynamics and other complex systems (Simonnet et
383 al., 2021; Hoffman et al., 2006; Weare, 2009; Vanden-Eijnden & Weare, 2013; Bouchet et al.,
384 2014; Chen et al., 2014; Farazmand & Sapsis, 2017; Dematteis et al., 2018; Mohamad & Sap-
385 s, 2018). A natural extension of these various techniques would combine elements of active rare
386 event sampling with the DGA method. Early developments of such a coupling procedure are pre-
387 sented in (Lucente et al., 2021).

6 Conclusion

Extreme weather events present a fundamental challenge to Earth system modeling. Many years of simulations are needed to generate sufficiently many extreme events to reduce statistical error, but high-fidelity models are needed to simulate those events accurately. Conventionally, no single model can provide both, simply because of computational costs. Here, we have demonstrated an alternative approach that leverages *ensembles of short*, high-fidelity weather model forecasts to calculate extreme weather statistics, with specific application to sudden stratospheric warming (SSW). By exploiting the huge database of forecasts stored in the subseasonal-to-seasonal (S2S) database (Vitart et al., 2017), we have obtained plausible estimates of the rate and seasonal distribution of SSW events that are (i) more precise, and (ii) more robust in distribution tails, than reanalysis data.

Our method uses data to estimate the dynamics on a subspace relevant for SSW, namely the polar vortex strength as measured by zonal-mean zonal wind. This single observable, augmented by time-delay embedding, gives a simple set of coordinates sufficient to estimate rate and seasonal distributions. Our demonstration opens the door to address many other data-limited questions of basic physical interest. For example, how important are vortex preconditioning and upward wave activity as triggers of SSW? (Charlton & Polvani, 2007; Albers & Birner, 2014). Do split-type and displacement-type events have fundamentally different mechanisms and/or different downstream effects? (Matthewman & Esler, 2011; Esler & Matthewman, 2011; O’Callaghan et al., 2014; Maycock & Hitchcock, 2015). Will climate change affect the frequency of SSW, perhaps through arctic amplification? (Charlton-Perez et al., 2008; Garfinkel et al., 2017; Kretschmer, Coumou, et al., 2018). How do other slow climatic variables, such as ENSO, the QBO, and the Aleutian Low affect SSW propensity? (Dimdore-Miles et al., 2021). These questions have been addressed in a number of coarse-resolution climate modeling studies, but high-resolution weather forecast data is an untapped source of potential for sharpening the answers. Our method offers a way forward, and is highly customizable to include physical features tailored for the problem at hand.

Another potential application of our methods is catastrophe modeling under climate change. Tropical cyclones pose a pressing problem for coastal communities, and have motivated several hybrid dynamical/statistical downscaling methods to project risk into the future under various climate change scenarios (Camargo et al., 2014; Lee et al., 2018; Jing & Lin, 2020; Sobel et al., 2021). Extreme precipitation of many varieties threatens cities and agriculture and is expected to change significantly with global warming (e.g., O’Gorman, 2012; Pfahl et al., 2017). Model resolution, again, is the limiting factor (Laflamme et al., 2016; O’Brien et al., 2016; He et al., 2019). Enlisting short weather forecasts, as we have done, may help identify precursors and drivers of changing frequency with unprecedented detail.

7 Open Research

Our analysis is based on publicly available datasets from the European Center for Medium-Range Weather Forecasts (ECMWF, 2022b) and from associated Copernicus Climate Data Store (ECMWF, 2022a). Python scripts to download the necessary data and reproduce the paper’s analysis will be made available in a public Zenodo repository at publication time.

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