

Downward migration of the zonal-mean circulation in the tropical atmosphere

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Key Points:

- Zonally coherent, or annular, fluctuations capture a higher fraction of circulation variability in the tropics than they do in midlatitudes
- Annular anomalies in tropical geopotential height and zonal wind migrate from the tropopause to the surface over approximately 10 days
- Downward migration is associated with a pulsing of the overturning Hadley circulation on subseasonal time scales

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Abstract

The annular modes of the extratropical atmosphere have received much attention for quantifying variability of the jet streams and storm tracks, despite the fact that the midlatitude circulation itself does not vary uniformly with longitude. While tropical fluctuations in geopotential height have lower amplitude than in the extratropics, they exhibit stronger zonal coherence, or dynamical annularity. A simple index is developed to characterize zonal-mean anomalies of the tropical circulation. It reveals that anomalies in geopotential height and zonal wind migrate downward from the upper troposphere to the surface on a time scale of about 10 days. These features are distinguishable from known modes of tropical variability, the Madden–Julian Oscillation in particular. Evidence from reanalysis and idealized model experiments confirms that this downward migration is quite generic and driven by mechanically forced variations in the strength of the Hadley circulation on subseasonal time scales.

Plain language summary: Earth’s atmosphere has certain recurring patterns which are “annular,” spanning an entire latitude circle. Annular patterns have proven to be surprisingly useful for weather and climate prediction. However, in the tropics, annular patterns have not been studied. We show that annular variability in the tropics does exist, and it contains interesting features that may be useful for prediction. In particular, zonal anomalies in the circulation migrate downward from the tropopause (about 16 km) to the surface over 10 days. These features are also apparent in a relatively simple climate model, which helps direct further research.

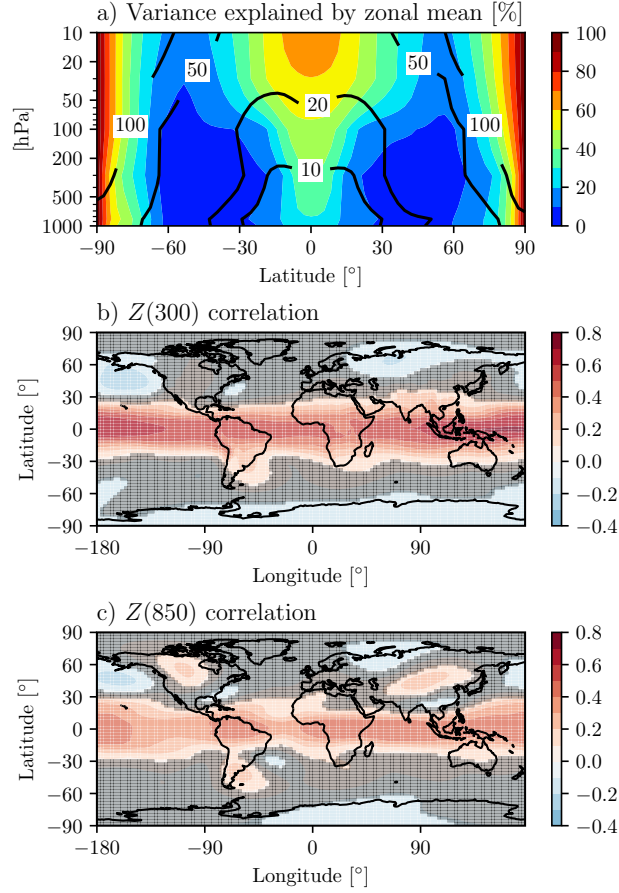
1 Dynamical annularity of the atmospheric circulation

The annular modes of the extratropical circulation have received much attention for their predictive utility and dynamical insight (see Kushner, 2010, for an extensive review). They characterize atmospheric variability in a remarkable variety of contexts, from unforced intraseasonal variability (e.g., Thompson & Wallace, 1998; Feldstein, 2000) to forced responses associated with global warming (Thompson et al., 2000) or stratospheric ozone loss and recovery (Thompson et al., 2002). The designation “annular” refers to their zonally uniform structure, which characterizes an exchange of mass and momentum between the polar cap and midlatitudes (Thompson & Wallace, 2000), although their dynamics fundamentally depends on zonally asymmetric eddies (Hartmann & Lo, 1998).

From the time of their discovery, however, there has been much debate on the extent to which an annular mode, or midlatitude variability more generally, is actually annular. Deser et al. (2000) and Ambaum et al. (2001) observed that while variability of geopotential height becomes zonally uniform over the polar cap in a trivial sense, as the latitude circle approaches the scale of the deformation radius, variations in the midlatitudes are not at all zonally uniform.

More recently, Gerber & Thompson (2017) attributed the zonal structure of the annular mode patterns to the “statistical annularity” of the midlatitude circulation. The annularity of the pattern reflects the fact that the statistics of geopotential height variability are fairly uniform across longitudes, as opposed to the fluctuations themselves. They observed, however, that geopotential height variations in the tropics, while of weaker amplitude, exhibit far more zonal coherence.

The zonal coherence—or “dynamical annularity”, to use the terminology of Gerber & Thompson (2017)—of geopotential height variability is quantified in Figure 1a. Here and throughout the study, data is taken from ERA–Interim reanalysis (Dee et al., 2011) from 1 January 1979 to 31 December 2017; JRA–55 (Kobayashi et al., 2015) reanalysis yields nearly identical results (not shown).



61 **Figure 1.** (a) The fraction of total variance associated with zonal mean anomalies in geopo-
 62 tential height (shaded), i.e., the power associated with wavenumber 0 in a zonal Fourier decom-
 63 position, normalized by the total power, and the RMS amplitude of the zonal mean flow (in
 64 meters; contoured). Unless specified otherwise, anomalies are defined as departures from a sea-
 65 sonally evolving climatology and high-pass filtered with a 1 year cutoff. (b) Correlation of 300
 66 hPa geopotential height anomalies with our annular variability index $\bar{Z}(300 \text{ hPa})$, described in
 67 the text. Hatching indicates regions where the correlation is not significantly different from 0
 68 with 95% confidence, assuming a decorrelation time of 40 days to determine the degrees of free-
 69 dom. (c) Same as (b), but showing the correlation of 850 hPa geopotential height anomalies with
 70 the $\bar{Z}(850 \text{ hPa})$ index.

71 In the midlatitude troposphere—from 30 to 60° latitude in both hemispheres—less
 72 than 10% of variability is characterized by annular fluctuations. Near the poles, the
 73 geometry of the sphere naturally leads to a dominance of zonally coherent motion, but an
 74 additional maximum is observed in the tropics. Zonally coherent fluctuations of
 75 geopotential height characterize thus approximately a quarter of tropical variability at
 76 lower levels, the fraction increasing to almost a half at the tropopause. In this study, we
 77 investigate the structure of this zonally coherent variability in the tropical troposphere.

2 Downward migration of annular circulation anomalies in the tropics

We first establish a convenient index to quantify annular variability in the tropics: zonal-mean geopotential height at the equator, computed on a given pressure level and denoted $\bar{Z}(p)$. We find that this simple index correlates very highly with more complex metrics, such as the leading time series from Principal Component Analysis of tropical geopotential height (cf. Baldwin & Thompson, 2009). Since these more complex options involve selecting parameters, we proceed with the simplest option. The data is filtered by subtracting the seasonal cycle and subsequently high-pass filtered with a 1 year cutoff to remove trends and low-frequency variability such as El Niño–Southern Oscillation (ENSO). As we discuss below, filtering is not critical.

Figure 1b,c shows the structure of variability associated with $\bar{Z}(850)$ and $\bar{Z}(300)$. The indices characterize broad variations across the entire tropics, extending approximately 20 degrees from the equator, particularly in the Pacific sector which dominates the zonal mean. Broad correlation in geopotential reflects the weak rotation in the tropical atmosphere. In the weak temperature gradient limit Sobel et al. (2001), any localized heating or cooling of the atmosphere (e.g., by convection) is balanced by ascent or descent, thereby homogenizing the temperature, and hence geopotential height. The enhanced positive correlation in Figure 1b relative to Figure 1c reflects the increase in annular variability with height shown in Figure 1a.

To explore the vertical coupling between annular anomalies in geopotential height, Figure 2a shows the cross-correlation of $\bar{Z}(p)$ with $\bar{Z}(850)$. We begin with unfiltered data (from which only the annual cycle has been removed), but even in this raw reanalysis data, we see a downward migration of geopotential height anomalies from the upper troposphere to the surface on a time scale of about 10 days, followed by the hint of an opposite signed anomaly which also migrates downward. A comparable picture emerges if we choose other levels (e.g., 300 hPa) as our base level of correlation.

Low-frequency variability (primarily ENSO) and trends lead to significant redness in the unfiltered cross-correlations. The subseasonal time scale becomes more apparent after applying a highpass filtering with 1 year cutoff (Figure 2b). This conservative approach to removing the red background of the unfiltered cross-correlations still allows in more than half of the variance at a time scale of 1 year, and reveals more clearly a downward-migrating signal, anticorrelated with preceding and proceeding anomalies. In testing we find that the results are not sensitive to the cutoff provided it is short enough to remove ENSO.

Several approaches to significance testing underscore the robustness of this phenomenon. Hatching in Figure 2 indicates statistical significance assuming a decorrelation time scale of 40 days, a conservative approach intended to under-estimate the effective degrees of freedom. Subsetting the dataset by decades leads to similar results, as does employing composites instead of cross-correlations (not shown). Significance will furthermore be assessed dynamically by comparing observations with numerical models.

We first ask: is this downward migration simply the zonal mean manifestation of a known tropical mode of variability? Both the MJO (Andersen & Kuang, 2012) and ENSO (Seager et al., 2003) are known to exhibit zonally coherent variability. With respect to the latter, the high pass filtering in Figure 2b is one approach to removing the ENSO signal. (Alternatively, removing ENSO variability by regressing out the signal linearly correlated with the Niño 3.4 index yields a very similar result to Figure 2a). The highpass filter, however, will tend to amplify the relative importance of the MJO. To ensure the downward migration exists independently of the MJO, we repeated the analysis after regressing away variability linearly correlated with the principal components of the

134 Real-time Multivariate MJO (RMM) index (Wheeler & Hendon, 2004). Again, the results
135 (not shown) are qualitatively identical with all MJO variability removed.

136 As an additional approach to answering this question, we consider a numerical
137 model which explicitly lacks ENSO or MJO-related variability. The Model of an Idealized
138 Moist Atmosphere (MiMA; Jucker & Gerber, 2017) is an idealized aquaplanet, and
139 modifications to the original model have been made to incorporate realistic zonal
140 asymmetries in the lower boundary; a later iteration of this configuration has been
141 published in Garfinkel et al. (2020). There is explicitly no oceanic variability and the
142 model does not capture the MJO. The model has no cloud feedbacks and only resolves the
143 large-scale circulation, employing a simplified parameterization of tropical convection.

144 Despite these substantial simplifications, geopotential anomalies also migrate
145 downward in MiMA, on comparable temporal and vertical scales (Figure 2c). If anything,
146 the cross-correlations in MiMA are slightly reduced in magnitude, implying that the full
147 physics of the actual atmosphere act to strengthen the downward migration, rather than
148 weaken it. Furthermore, the horizontal structure of tropical geopotential height variations
149 in the model is comparable to Figure 1b,c (not shown).

150 We have shown idealized moist aquaplanet integrations because this is perhaps the
151 simplest configuration that can reproduce the phenomenon. When one removes the
152 impact of moisture from the model, as in the idealized atmospheric model of Held &
153 Suarez (1994), the downward migration of geopotential anomalies is not apparent, as we
154 will address in Section 4. Dry dynamical cores are fully capable of capturing the annular
155 modes (e.g., Gerber & Vallis, 2007), or the more recently discovered Baroclinic Annular
156 Modes (Barnes & Thompson, 2014), but our results indicate that additional complexity is
157 critical to capturing the annular variations of the tropics.

158 Although the MJO is not necessary for the downward-migrating signal in equatorial
159 geopotential height, it does exhibit a similar signal. To quantify the MJO’s signature on
160 zonal-mean geopotential, we construct weighted composites of normalized geopotential
161 anomalies at each level as a function of RMM phase (Wheeler & Hendon, 2004). Noting
162 that 8 phases of the oscillation correspond to approximately 40 days, these composites
163 show qualitatively similar downward migration on consistent time scales (Figure 2d).
164 Thus in isolation an MJO event is sufficient, but not necessary, to achieve downward
165 migration.

166 To get a sense of the actual variability, an example of downward migration for both
167 positive and negative anomalies is shown in Figure 2e. Their vertical coherence and time
168 scales are similar to the cross-correlated picture, and a lagged anticorrelation is apparent.
169 These events and the cross-correlated average are reminiscent of the “dripping paint
170 plots” of extratropical annular mode composites (Baldwin & Dunkerton, 2001), although
171 the resemblance is superficial, as the dynamics of the two are quite distinct.

172 Downward migration in the Northern and Southern annular mode indices involves
173 coupling between the tropospheric extratropical jet and the stratospheric polar vortex: a
174 deceleration of the polar vortex is usually followed by a nearly barotropic response in the
175 troposphere (Thompson & Wallace, 2000; Baldwin & Dunkerton, 2001). In the tropics,
176 the downward migration is entirely within the troposphere. In addition, downward
177 migration in the extratropical annular modes, while very robust, only appears in
178 composite analysis and is not evident from simple cross-correlation, unlike the tropics.

179 Downward migration in the tropical circulation is also apparent in other
180 meteorological variables, as shown in cross correlation of the zonal mean zonal wind with
181 $\bar{Z}(850)$ in Figure 3a. Weakened trade winds migrate downward in association with
182 anomalously high geopotential height, reflecting cyclostrophic balance. Conversely,
183 strengthened trade wind anomalies migrate downward in concert with negative
184 geopotential height anomalies. In the next section, we parse the structure of downward

185 migrating geopotential and zonal wind by cross-correlating with the key dynamical fields
 186 at play.

187 **3 Mechanically forced variations in the Hadley Circulation**

188 Geopotential height is the altitude of a given pressure surface of the atmosphere,
 189 and thus depends on both the surface pressure and the temperature of the atmosphere
 190 below. A pressure surface is elevated when the surface pressure increases; that is, when
 191 the local mass of the atmospheric column is increased, requiring one to ascend higher
 192 through the atmosphere to reach a given pressure. Similarly, a pressure surface is elevated
 193 if the atmosphere below it is warmer, and so less dense, again requiring a higher ascent to
 194 reach a given pressure. The downward migration of geopotential height anomalies shown
 195 in Figure 2 thus corresponds to a coordination, or phase locking, between temperature
 196 anomalies and surface pressure anomalies in the tropics.

204 Figure 3b,c show that anomalously high low-level geopotential (or surface pressure)
 205 is preceded by anomalously warm temperatures throughout the tropical troposphere,
 206 which lifts geopotential height at upper levels. Correlation between low-level geopotential
 207 and temperature reaches a maximum at a lead of 10 days (-10 on the horizontal axis), at
 208 which time surface pressure anomalies are uncorrelated: the atmosphere is warmest when
 209 surface pressure is neutral. The atmosphere cools as surface pressure anomalies build,
 210 with temperature anomalies vanishing at lag 0, just as surface pressure reaches a
 211 maximum. The atmosphere continues to cool, as surface pressure relaxes, reaching a
 212 maximum negative anomaly near a lag of 10 days (10 on the horizontal axis), maintaining
 213 quadrature with the surface pressure field.

214 Since the surface pressure indicates the total mass of the atmospheric column,
 215 downward migration in tropical geopotential can thus be described as a progression from
 216 *hot to heavy* (more mass) to *cold*. The cross-correlation is symmetric, by construction, but
 217 additional testing via composite analysis (not shown) reveals that this symmetry does
 218 hold between positive and negative anomalies: downward migration of negative
 219 geopotential height anomalies are associated with a cold–light–warm pattern.

220 This phase locking between temperature and surface pressure reflects a common
 221 dynamical origin: a pulsing of the Hadley circulation, illustrated in Figure 3d. On
 222 subseasonal time scales, the mean overturning circulation of the tropics is dominated by
 223 localized convection and associated downdrafts: hence only a weak, albeit statistically
 224 significant, zonal mean signal emerges from a noisy background. The anomalous warming
 225 of the tropics peaking at -10 days (Figure 3b) is associated with a suppression of
 226 convection (indicated by reduced precipitation, and consistent with an increase in gross
 227 moist stability) and a *weakening* in tropical upwelling (indicated by the positive
 228 correlation with the pressure velocity ω at 300 hPa). Upwelling is reduced at all
 229 tropospheric levels, suggesting an overall slowing of the overturning.

230 The anticorrelation between temperature and the upwelling velocity suggests that
 231 these anomalies are mechanically forced. Indeed, to generate kinetic energy
 232 thermodynamically, warm air must be associated with *increased* upwelling. The positive
 233 temperature anomalies (or equivalently, increase in moist stability) are consistent with a
 234 reduction in cooling driven by adiabatic expansion. Likewise, the reduction in upwelling
 235 throughout the tropical atmosphere is associated with a buildup of mass, such that
 236 surface pressure continues to increase until the overturning circulation begins to speed up
 237 again at lag 0.

238 The synchronous variation between ω and temperature would be expected on time
 239 scales where the radiative damping is sufficiently fast to bring the atmosphere in balance
 240 with mechanically forced upwelling anomalies. The effective radiative damping time scale
 241 of the tropical atmosphere at lower levels is sufficiently fast relative to the subseasonal

242 time scale of these oscillations (Jucker et al., 2013), but this is not the case at upper
 243 levels. Since precipitation is suppressed (Figure 3d) we conclude that latent heating is not
 244 responsible for warming at upper levels, but further work is needed to determine how the
 245 upper troposphere temperature anomalies remain synchronized.

246 4 Downward migration in idealized atmospheric models

247 As a probe of the downward migrating mechanism, we turn back to atmospheric
 248 model experiments. As shown in Figure 2c, downward migration was observed in MiMA
 249 which explicitly neglects clouds processes (see Jucker & Gerber, 2017, for further details).
 250 Thus cloud radiative feedbacks are unlikely to play a critical role in the mechanism.
 251 However, MiMA includes an idealized representation of land surfaces and oceanic heat
 252 transport to generate realistic stationary waves (as described in Garfinkel et al., 2020).
 253 Downward migration is nearly absent, however, in aquaplanet configurations in which all
 254 sources of zonal asymmetry in the lower boundary were removed. Hence stationary waves,
 255 or at least their impact on mean climatology and variability, are important.

256 Given the evidence in the previous section for a mechanical (momentum) forcing in
 257 the mechanism, we next probe the phenomenon in a dry atmospheric model. Here,
 258 radiative and moist processes are explicitly removed, and all thermodynamical processes
 259 are approximated by Newtonian relaxation to an equilibrium solstice temperature profile
 260 (Held & Suarez, 1994; Polvani & Kushner, 2002). The model nonetheless produces a fairly
 261 realistic Hadley Circulation, and vigorous midlatitude variability drives variations in
 262 tropical geopotential.

263 As seen in Figure 4a, the simplified physics of the dynamical core, however, are
 264 insufficient to generate a downward migration of geopotential anomalies. Use of an
 265 equilibrium equinox temperature profile, or integrations with large scale topography
 266 (Gerber & Polvani, 2009), not shown, exhibit a similar absence of downward migration.
 267 This indicates that momentum from the extratropics, associated with Rossby waves that
 268 break deep in the tropics, is insufficient to drive variations in the Hadley Cell that can
 269 phase lock the temperature and surface pressure variability.

270 While the moist and dry models share similar midlatitude variability, the moist
 271 model captures a more realistic representation of tropical variability, in particular by
 272 simulating convectively coupled waves. The model is an extension of that explored by
 273 Frierson (2007) and shares a similar spectrum of tropical waves. To test whether the
 274 excitation of tropical waves are important, we introduce a simplistic tropical variability
 275 into the dry dynamical core through a prescribed temperature tendency.

276 Figure 4b shows the impact of artificially excited periodic variations in the tropical
 277 atmosphere, given by

$$278 \quad \frac{\partial T}{\partial t} = \begin{cases} A \cos\left(2\pi \frac{t}{\tau}\right) \cos\left(\frac{\pi}{2} \frac{\phi}{\phi_0}\right) & |\phi| < \phi_0 \\ 0 & |\phi| > \phi_0 \end{cases} \quad (1)$$

279 where $A = 0.03$ K/day, $\tau = 32$ days, and $\phi_0 = 20$ degrees, chosen to mimic the
 280 features seen in reanalysis. The temperature perturbations are zonally and vertically
 281 uniform, since variations in the thermal damping time scale act to recover approximately
 282 the vertical structure seen in Figure 3a. When this periodic temperature forcing is applied
 283 in the dynamical core, downward migration re-emerges, reflecting the previous phase
 284 locking between temperature and surface pressure seen in reanalysis and MiMA. For these
 285 integrations, highpass filtering is applied at a cutoff of 90 days instead of 1 year, to
 286 compensate for its lack of seasonal cycle.

287 Experimentation with the vertical, meridional, and temporal variation of the
 288 temperature tendency indicates that the downward migration is fairly generic provided
 289 the time scale of forcing remains subseasonal, and the amplitude is strong enough to
 290 excite anomalies on the order of a few Kelvin. This suggests that tropically forced
 291 variability, associated with convectively coupled waves, may play a crucial role in the
 292 mechanism.

297 5 Summary

298 Annular variations in geopotential height capture a higher fraction of the natural
 299 variability in the tropical atmosphere as compared to the midlatitudes. We constructed a
 300 simple index to characterize the annular variability of the tropics: zonal-mean
 301 geopotential at the equator. The index tracks broad variations of the height fields at each
 302 level, extending approximately 20 degrees in latitude into both hemispheres.
 303 Cross-correlation between levels reveals downward-migrating circulation anomalies on
 304 intraseasonal time scales. Positive geopotential height anomalies, associated with
 305 decreased trade winds, migrate from the tropopause to the surface on a time scale of
 306 approximately 10 days. Likewise, enhanced trade winds migrate downward with negative
 307 geopotential height anomalies.

308 While the Madden–Julian Oscillation (MJO) is associated with a similar zonal mean
 309 signature, we find that these downward migration circulation anomalies are more generic.
 310 They reflect a phase locking of temperature and surface pressure anomalies associated
 311 with variations in the overturning circulation on intraseasonal time scales. Anticorrelation
 312 between the vertical velocity and temperature suggests that these variations are
 313 mechanically driven. Simulations with idealized atmospheric models attest to the generic
 314 nature of these downward migrating anomalies, but indicate that the mechanism depends
 315 on the presence of tropical variability. Future work intends to explore the mechanism in
 316 greater detail, and discuss the predictability of these oscillations.

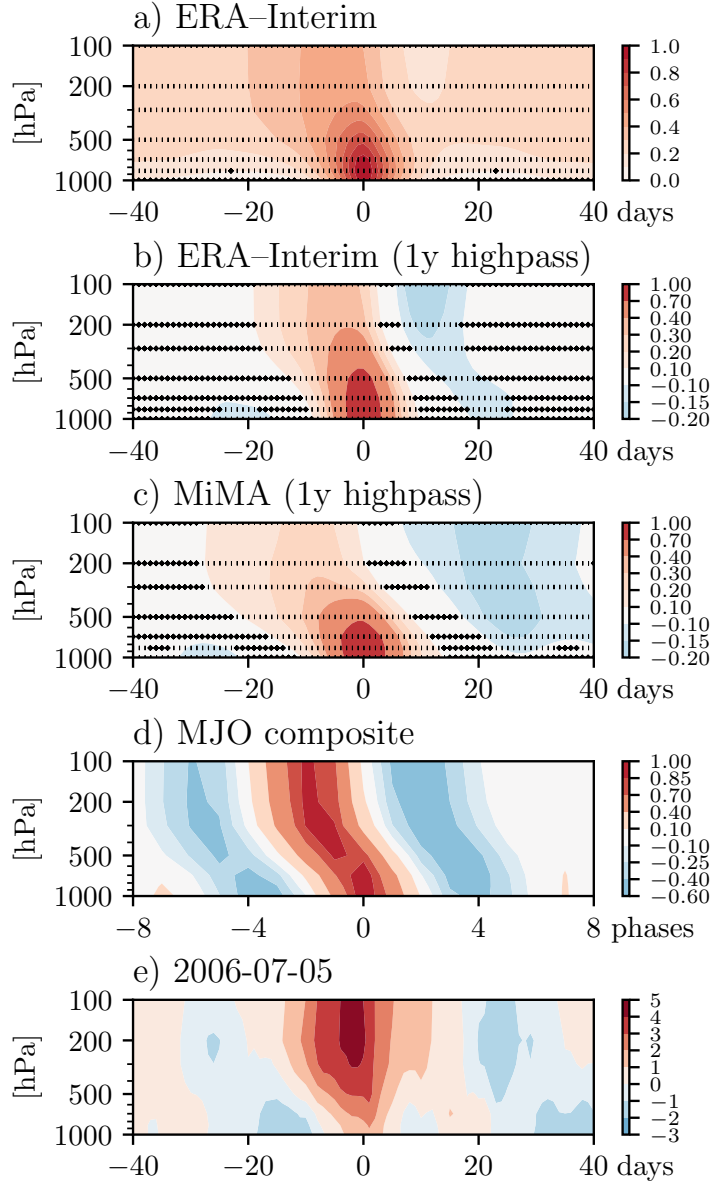
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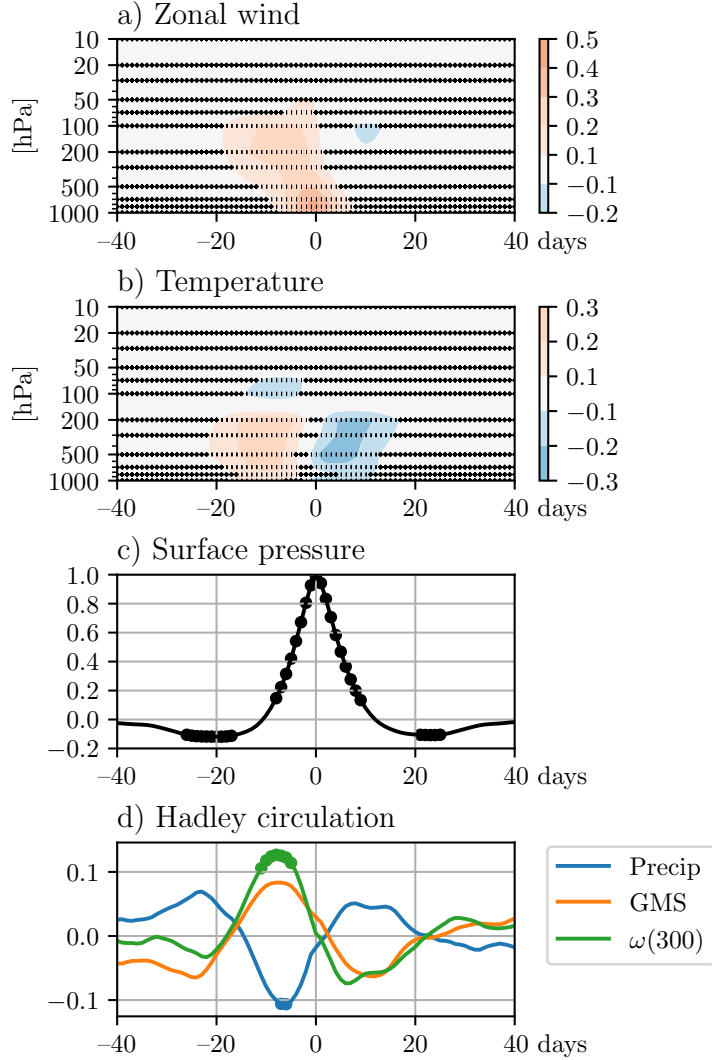
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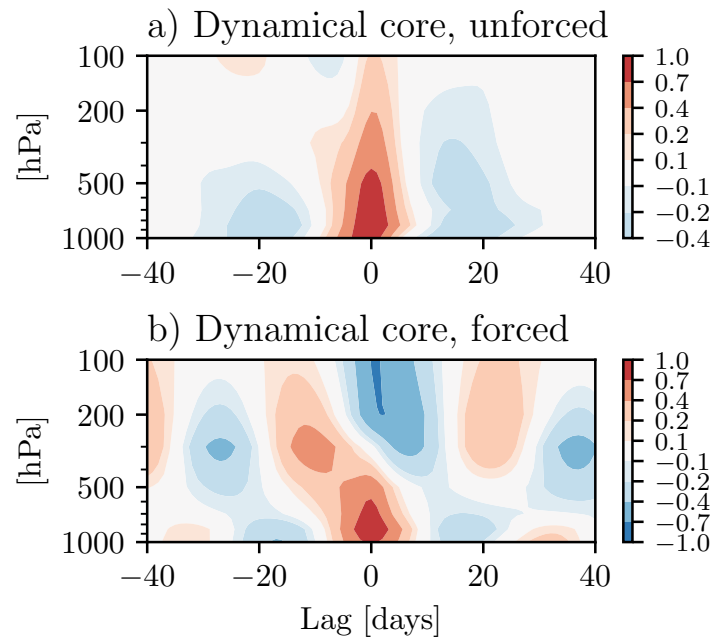
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104 **Figure 2.** (a–d) Cross-correlation of $\bar{Z}(p)$ with $\bar{Z}(850)$. Data (a) from ERA–Interim unfil-
 105 tered, (b) from ERA–Interim with a 1-year highpass filter, (c) from MiMA with a 1-year highpass
 106 filter, (d) composited by MJO RMM phase (amplitude-weighted; see text for details). (e) ERA–
 107 Interim $\bar{Z}(p)$ with a 1-year highpass filter, normalized to have unit variance, centered around July
 108 5, 2006. In (a–c), hatching indicates inability to reject the null hypothesis that the correlation is
 109 zero with 95% confidence, using a decorrelation time of 40 days.



197 **Figure 3.** Cross-correlation of $\bar{Z}(850)$ with anomalous zonally averaged equatorial (a) tem-
 198 perature, (b) zonal wind, (c) surface pressure p_s , and (d) Hadley circulation metrics. In (d), the
 199 Hadley circulation is characterized by precipitation rate, gross moist stability (GMS; moist static
 200 energy at 850 hPa minus 200 hPa), and vertical velocity ω at 300 hPa. Data from ERA-Interim
 201 is, as before, deseasonalized and 1-year highpass filtered. Hatching (a,b) and missing dots (c,d)
 202 indicate inability to reject the null hypothesis at 95% confidence, using a decorrelation time of 40
 203 days.



293 **Figure 4.** Cross-correlation of $\bar{Z}(p)$ with $\bar{Z}(850)$ in a dry dynamical core with Held & Suarez
 294 (1994) relaxation physics in (a) a control integration and (b) forced with the temperature ten-
 295 dency described in equation (1). Here the 1-year highpass filter is replaced by a 90-day bandpass
 296 filter, to compensate for the lack of seasonal cycle in the model.